

The Standard Model Higgs Search at the Tevatron

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On behalf of the DØ Collaboration

On behalf of the Tevatron Higgs Working Group

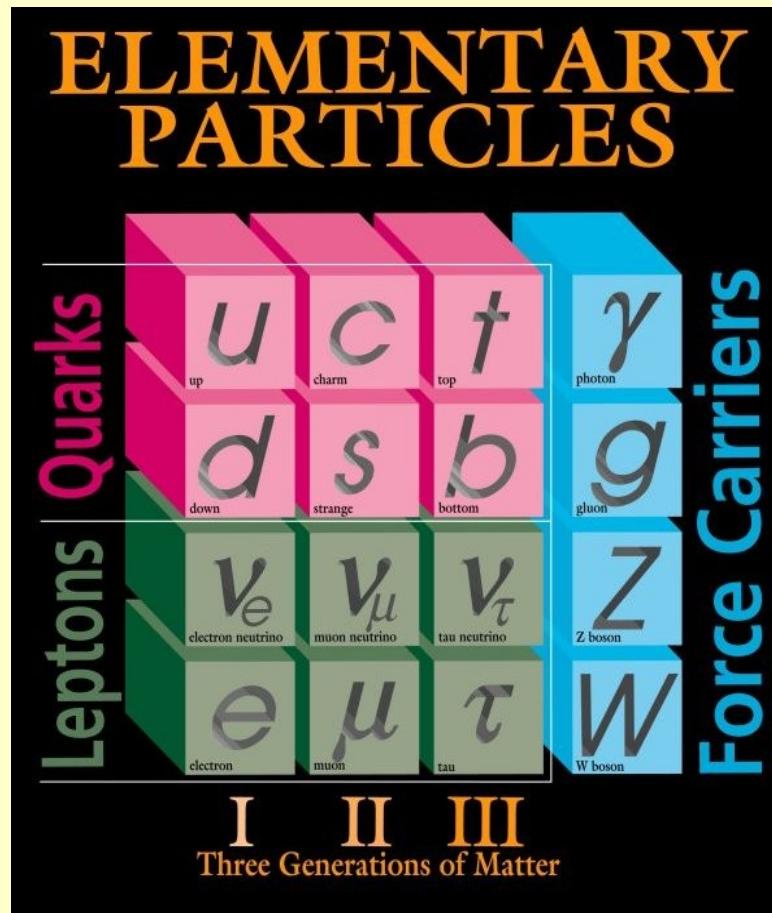


The Case for the Higgs



MSU HEP Seminar
Sept 25th, 2007

- Many years of work have led to our current description of matter and its interactions: ***the Standard Model***



Cast of characters includes

Matter particles (fermions): quarks and leptons

Force carriers (bosons): photon, gluon, W^\pm/Z^0

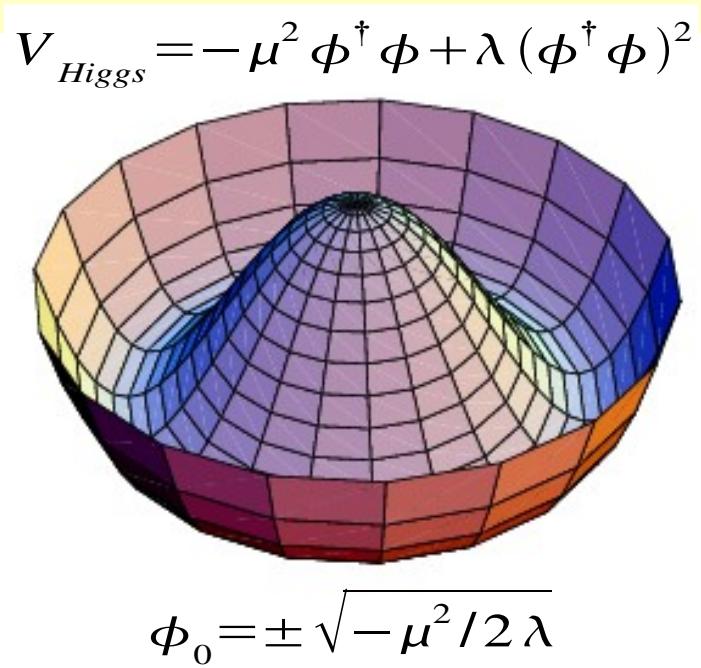
Highly successful predictive model

But there's a problem!! No explanation for particle masses

The Case for the Higgs

- ✗ Electroweak model is very powerful
 - $SU(2)_L \times U(1)_Y$ well-tested at collider experiments

But it is not a symmetry of our vacuum – otherwise gauge bosons would be massless
- ✗ Higgs mechanism provides a natural solution
 - Add one complex doublet of scalar fields in a Φ^4 potential
- ✗ Symmetric solution unstable: broken EW symmetry creates non-zero vacuum expectation value
 - W^\pm/Z^0 longitudinal polarizations absorb three degrees of freedom, remaining one becomes neutral scalar (Higgs boson)
 - Ground state VEV parameterizes W/Z masses
 - Higgs mass not predicted: $m_H \propto \mu$*



An Aside:

Quark and lepton masses must be added “by hand” via Yukawa-type terms in the Lagrangian. The Higgs mechanism is convenient, but may not be the complete picture.

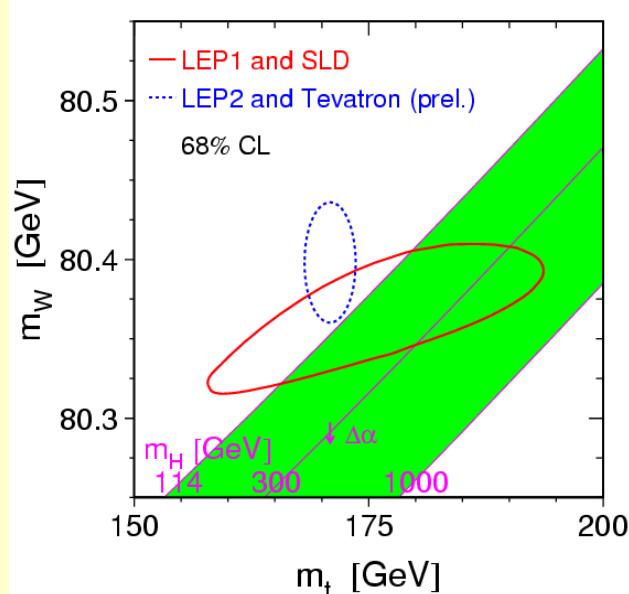
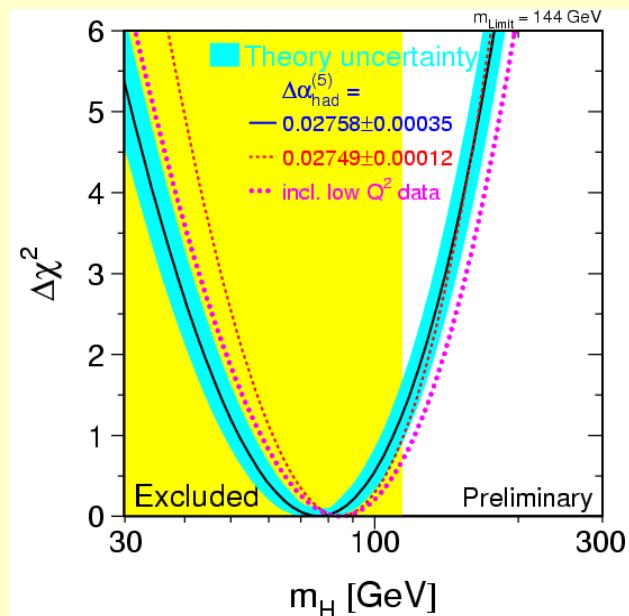
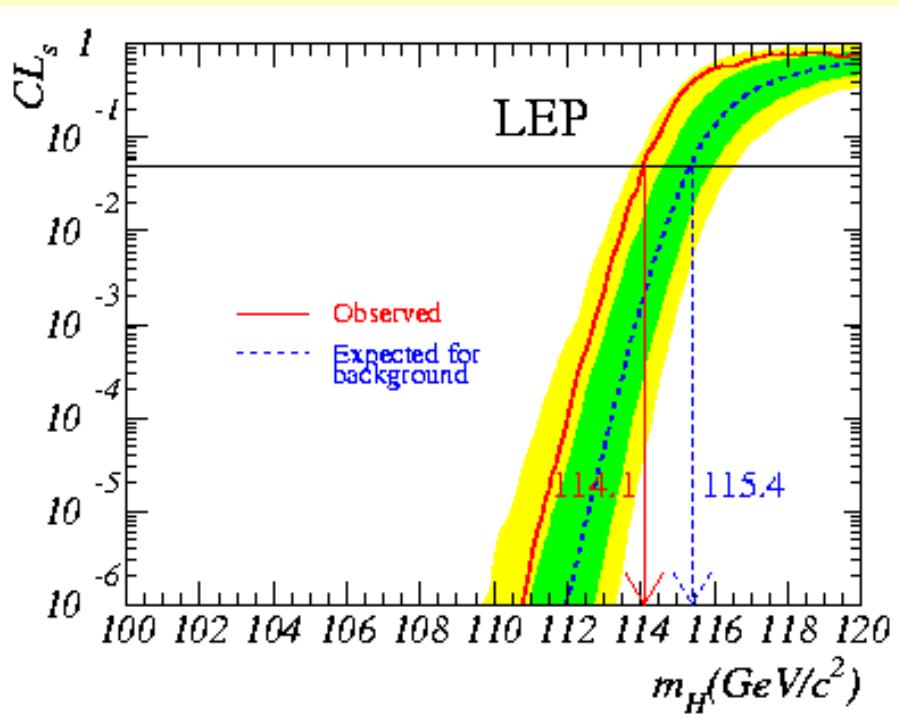
Evidence for the Higgs

- Direct searches and indirect evidence have restricted possible phase space

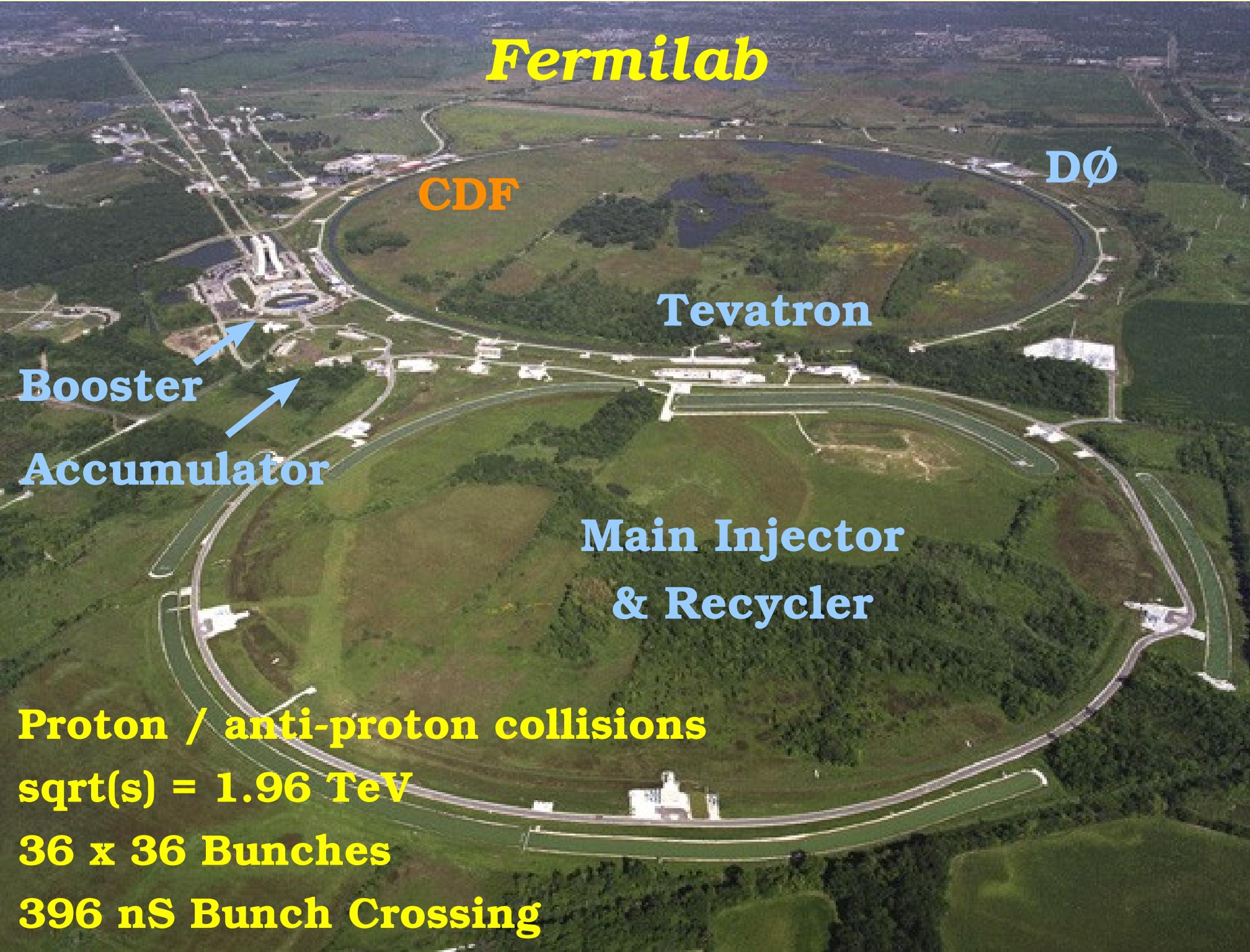
LEP II direct search: $m_H > 114.4$ GeV (95% CL)

Continuous refinements of top-quark and W-boson masses at the Tevatron

Fit of precision electroweak data: $m_H < 144$ GeV at 95% CL ($m_H < 182$ GeV including LEP II limit)



Fermilab



CDF

DØ

Tevatron

Main Injector
& Recycler

Booster
Accumulator

Proton / anti-proton collisions

$\text{sqrt}(s) = 1.96 \text{ TeV}$

36 x 36 Bunches

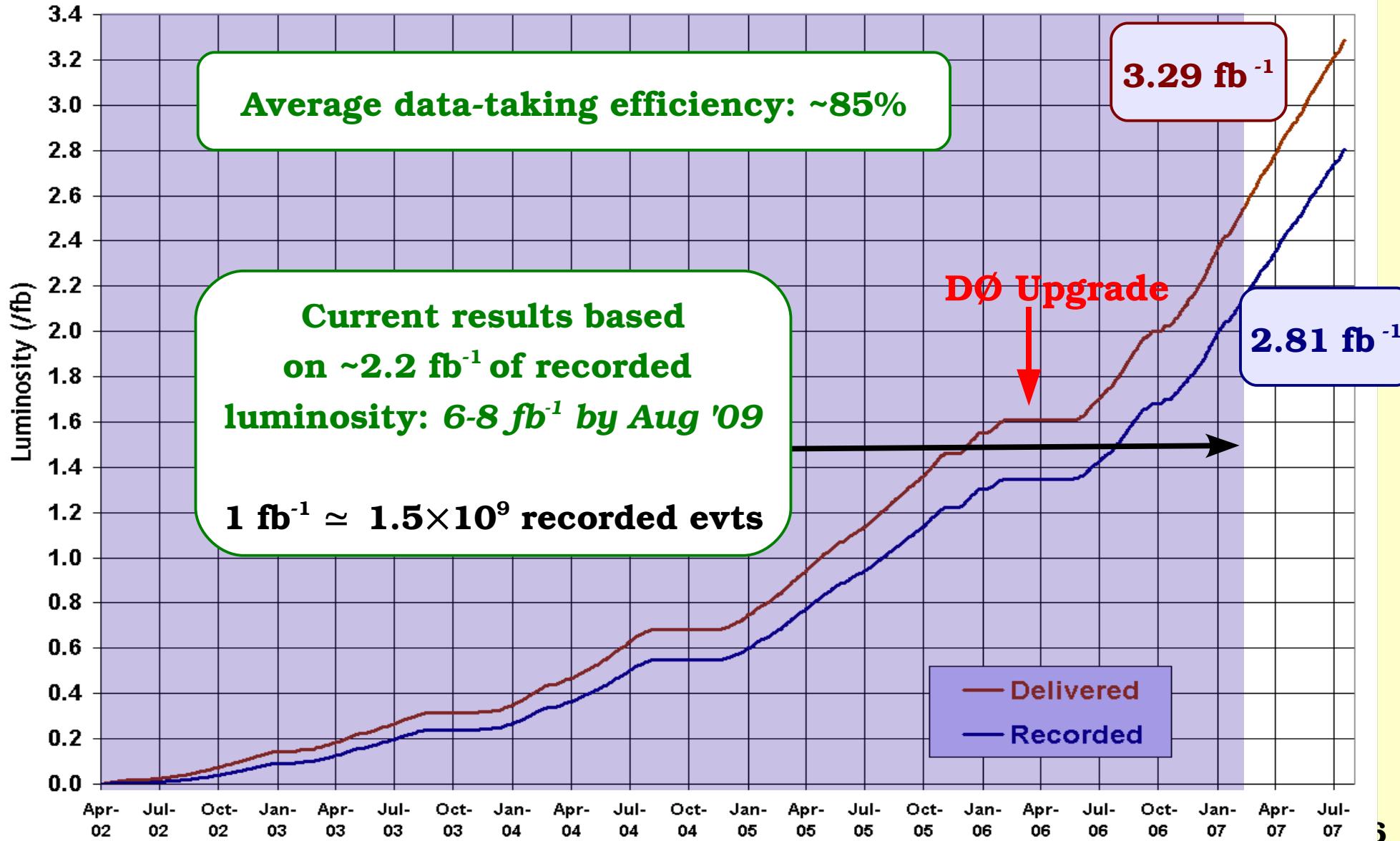
396 nS Bunch Crossing

Dataset

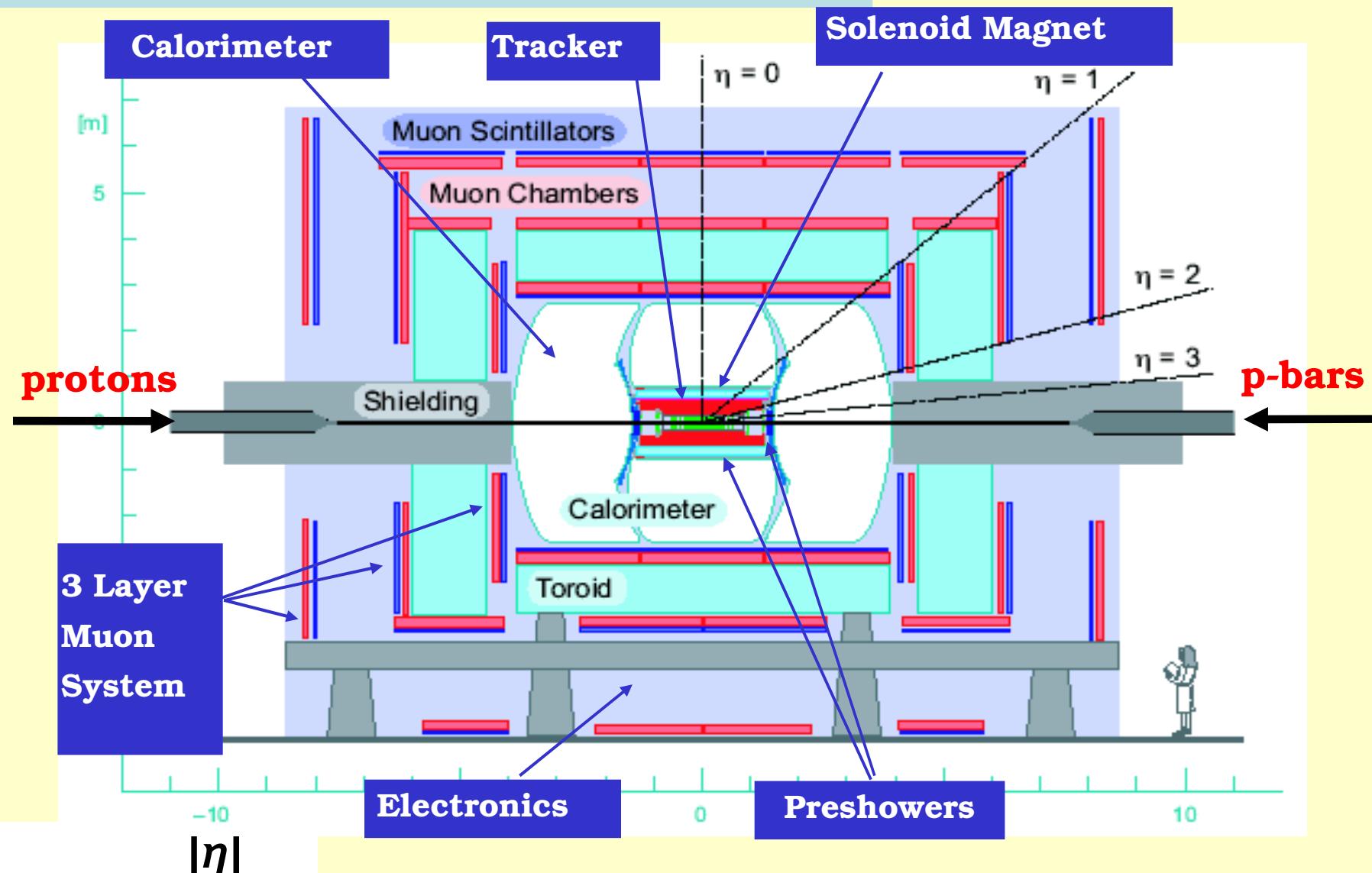


Run II Integrated Luminosity

19 April 2002 - 5 August 2007



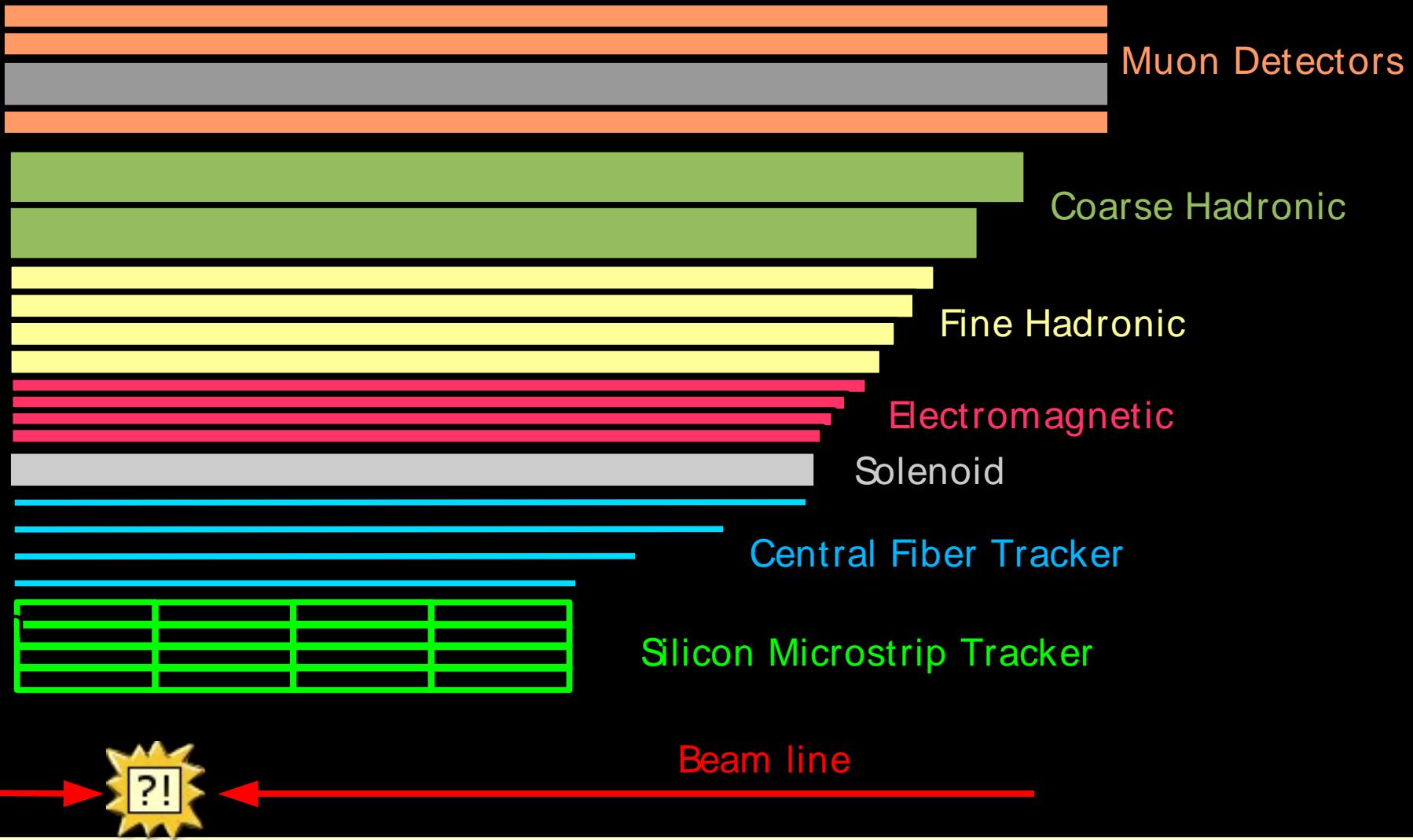
The DØ Detector



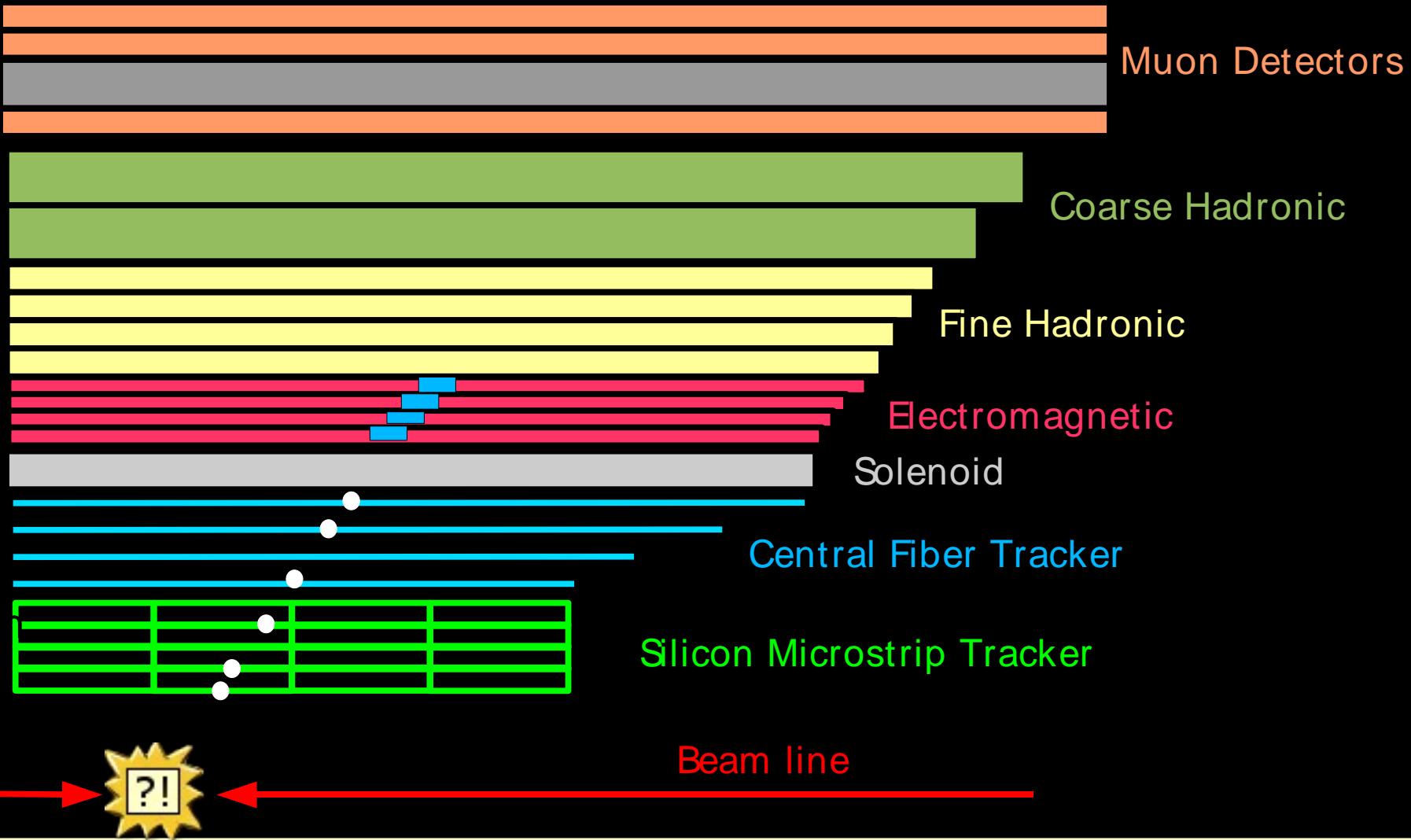
	$ \eta $
Muon ID	~ 2
Tracking	~ 2.5
EM / Jets	~ 4

The CDF detector has similar features of detector design

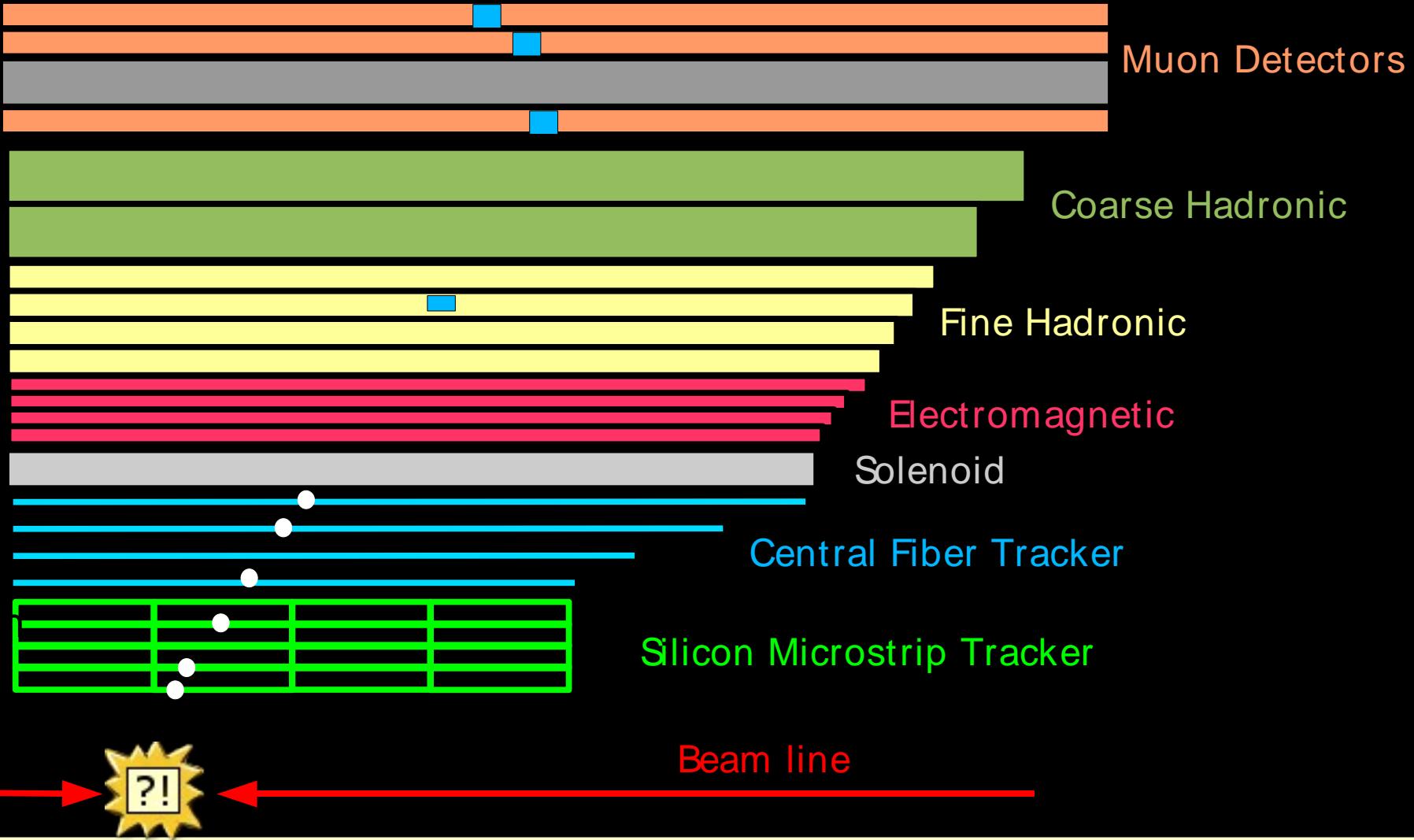
Particle Detection



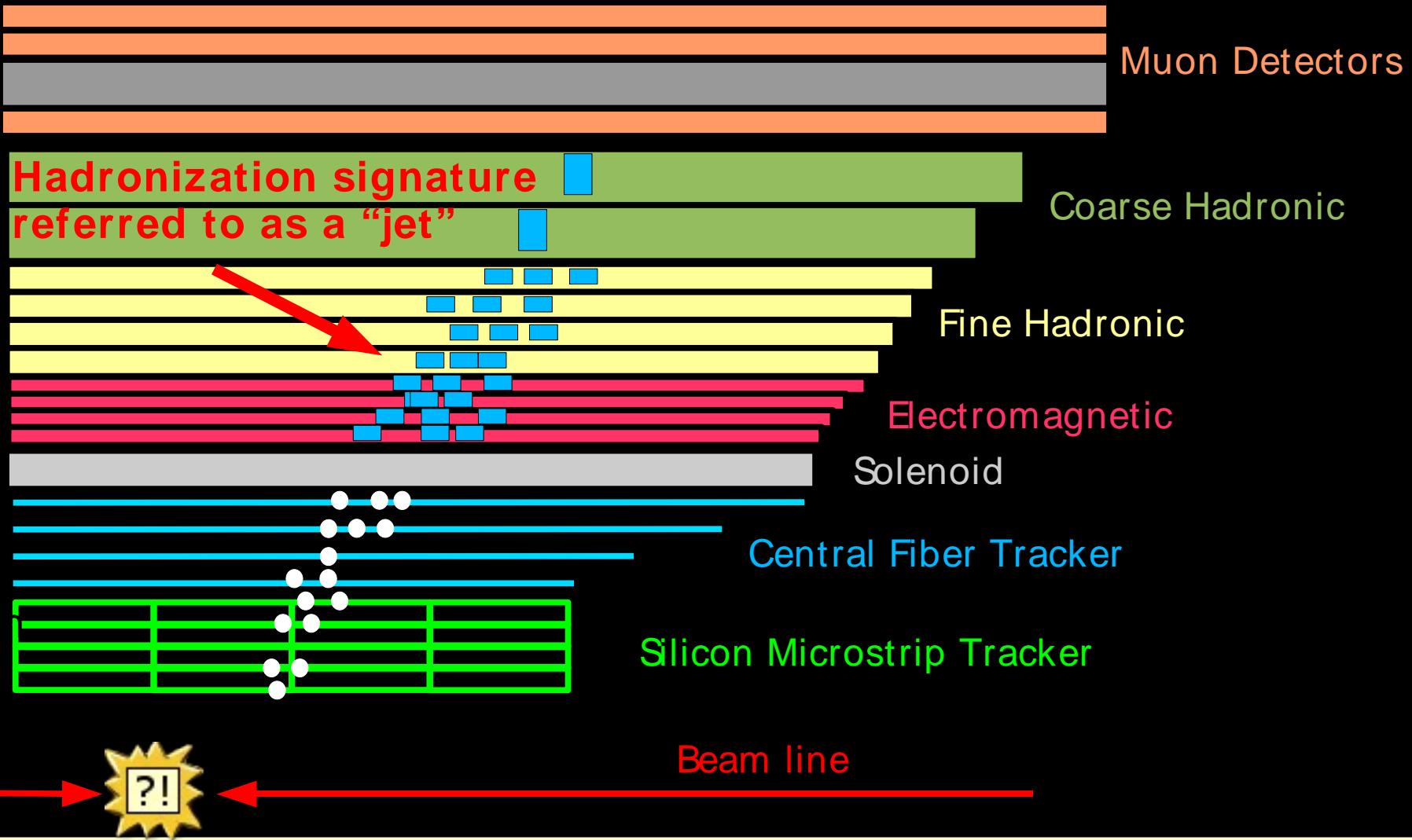
Particle Detection: Electrons



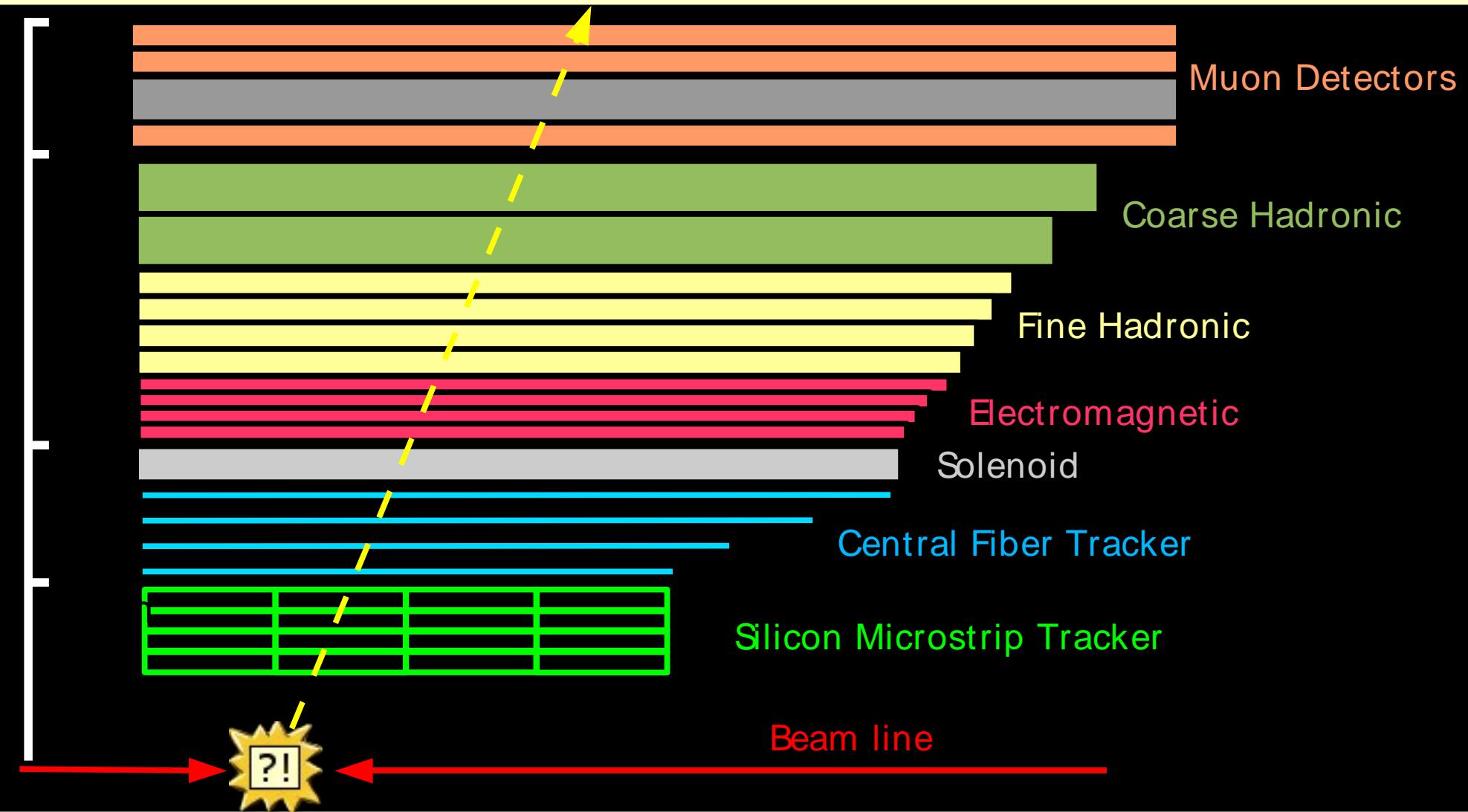
Particle Detection: Muons



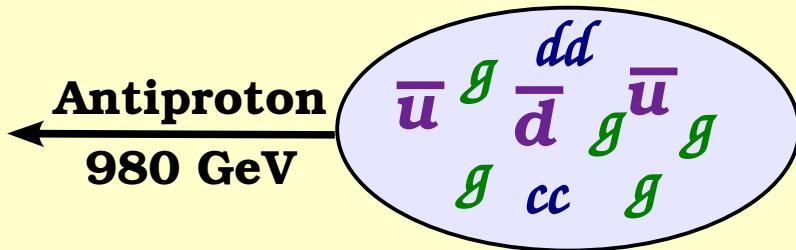
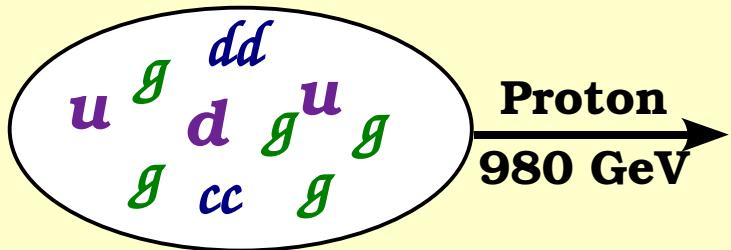
Particle Detection: Quarks & Gluons



Particle Detection: Neutrinos

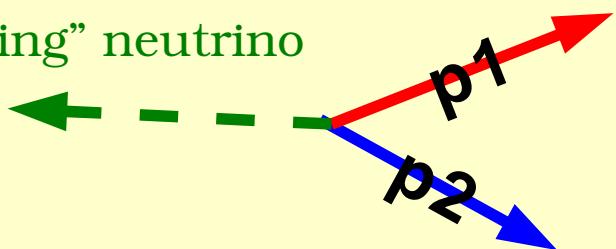


Proton-Antiproton Collisions



Transverse Variables

“missing” neutrino



$$MET = \sqrt{\left| \sum \vec{P}_x \right|^2 + \left| \sum \vec{P}_y \right|^2}$$

- ✗ Summation performed over calorimeter & muon energies
- Signature of non-interacting neutrinos
- ✗ Capitalize on momentum balance in transverse plane
- “Transverse momentum”, pT

Coordinate Systems

- ✗ Based on relativistic invariance
 - Allows for consistent description of z-boosted events
- ✗ Azimuthal angle ϕ is manifestly invariant under boosts, but θ isn't.
- ✗ Introduce invariant “rapidity” and zero-mass limit “pseudo-rapidity”

$$\text{Rapidity} = \frac{1}{2} \frac{\ln(E + P_z)}{\ln(E - P_z)}$$

$$\eta = -\ln \left(\tan \left(\frac{\theta}{2} \right) \right)$$

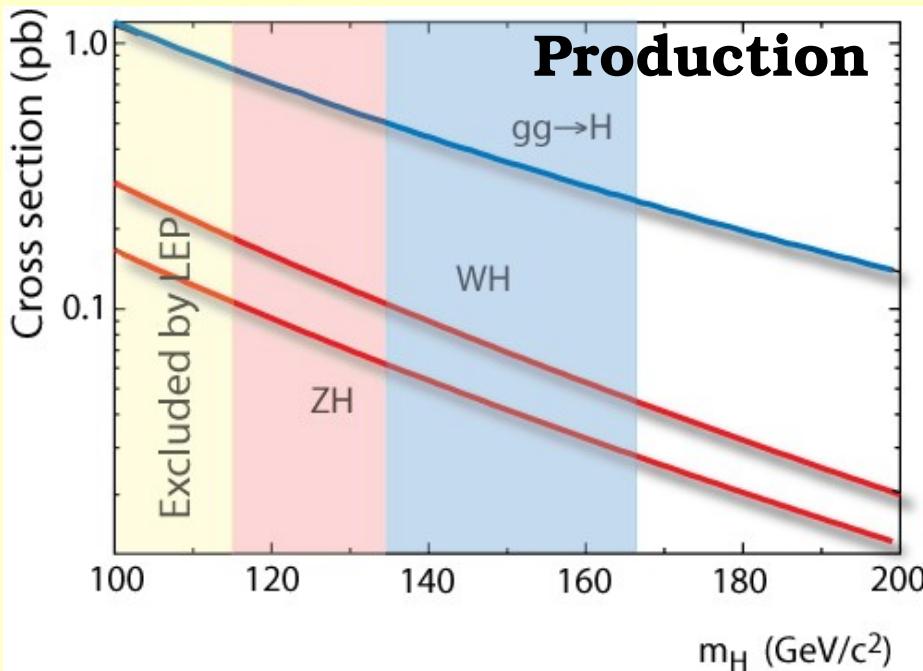
SM Higgs at the Tevatron

Gluon fusion dominant at hadron colliders

Large backgrounds restrict many useful Higgs decay channels

Next largest is associated production of W/Z + Higgs

Leptonic decays of W/Z bosons provide tag for trigger and analysis

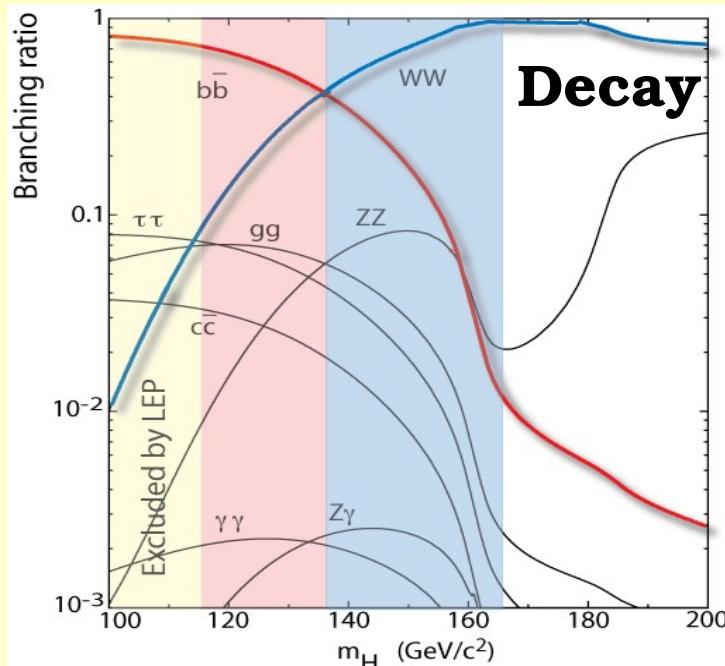


Low-mass Higgs ($m_H < 135$ GeV) prefers to decay to bottom-quark pairs

Need efficient identification of bottom quarks to reduce backgrounds

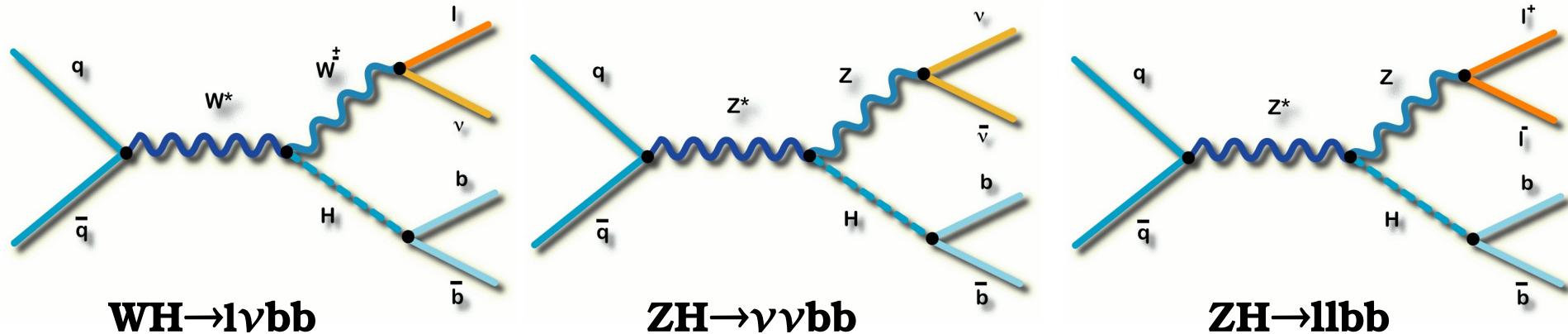
At high mass ($m_H > 135$ GeV), search for $H \rightarrow WW^*$

Off-shell W boson allows off-resonance production



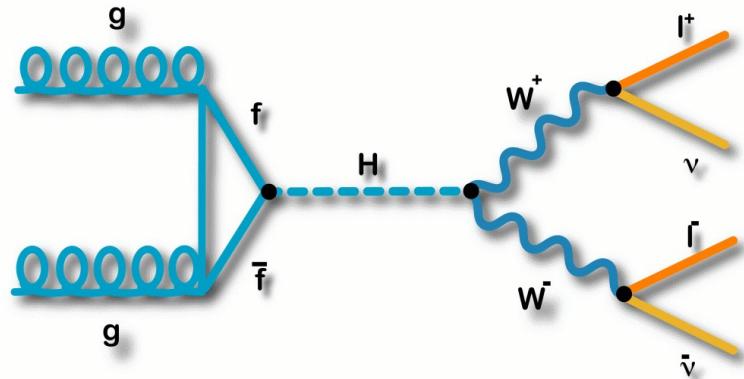
Search Channels

Associated Production: Low mass only, 3 final states



Gluon Fusion Production:

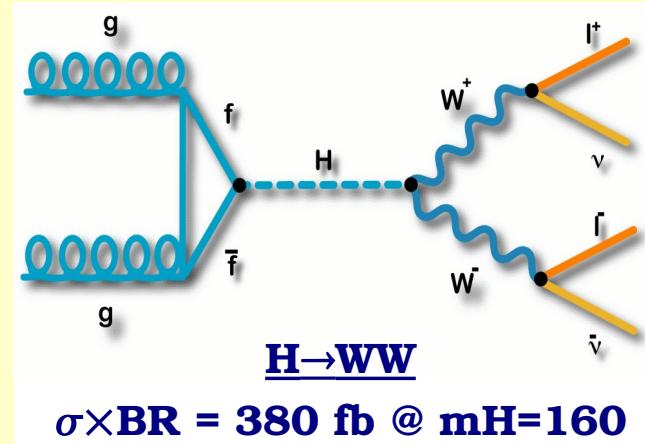
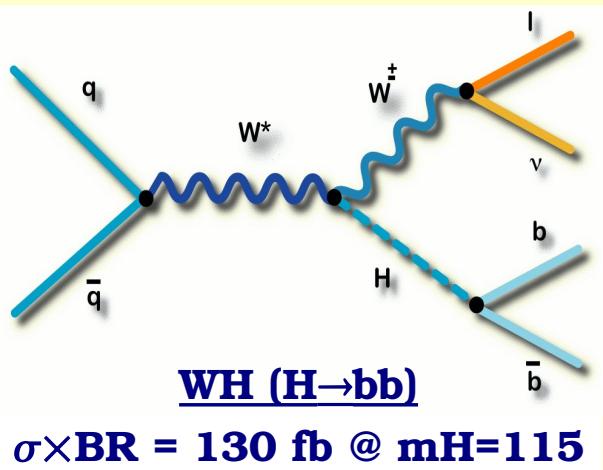
Maximum sensitivity at high mass,
also useful at low mass



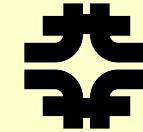
Backgrounds & Rates



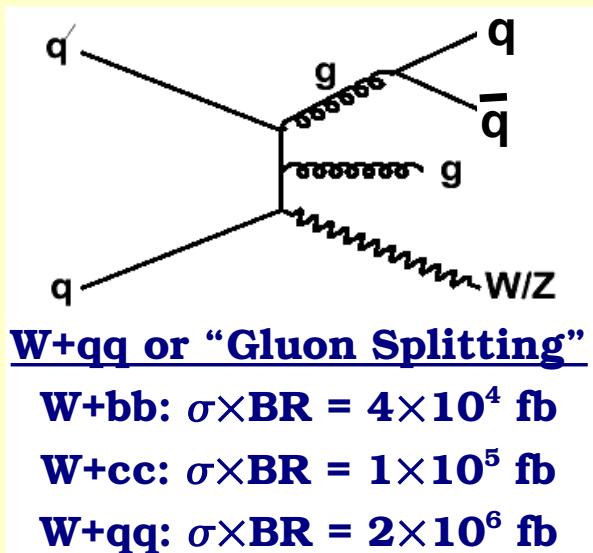
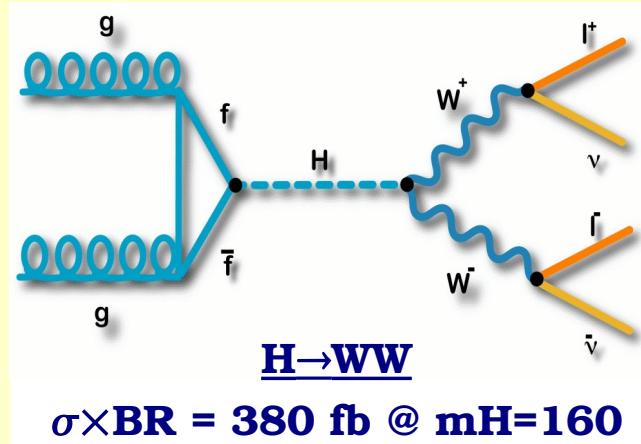
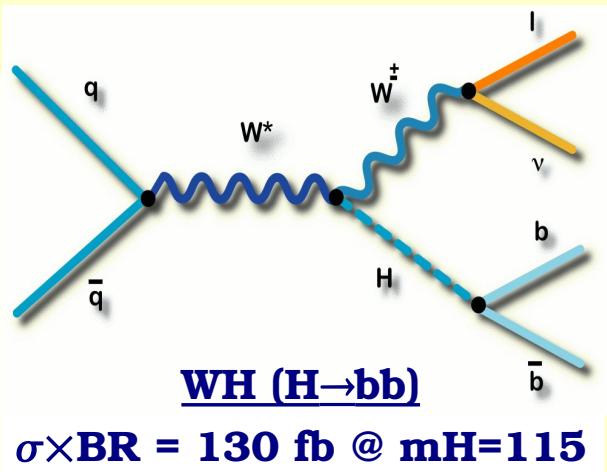
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Backgrounds & Rates



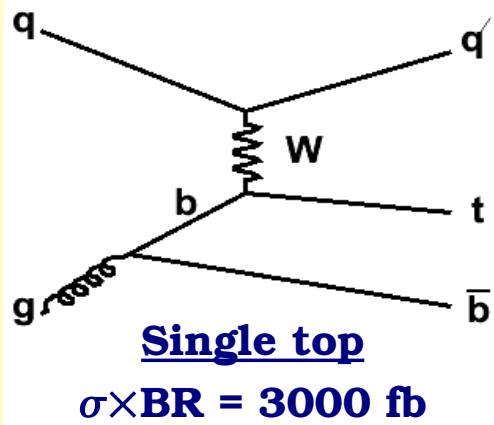
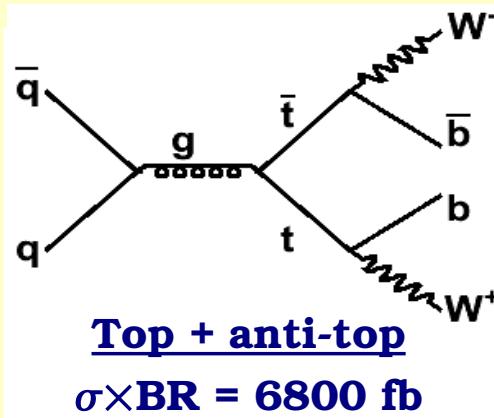
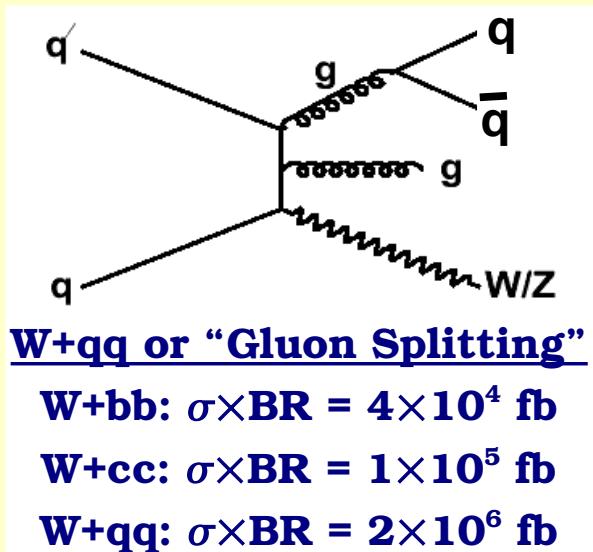
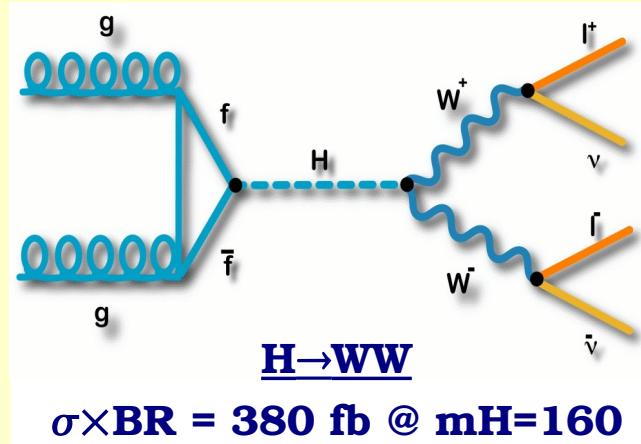
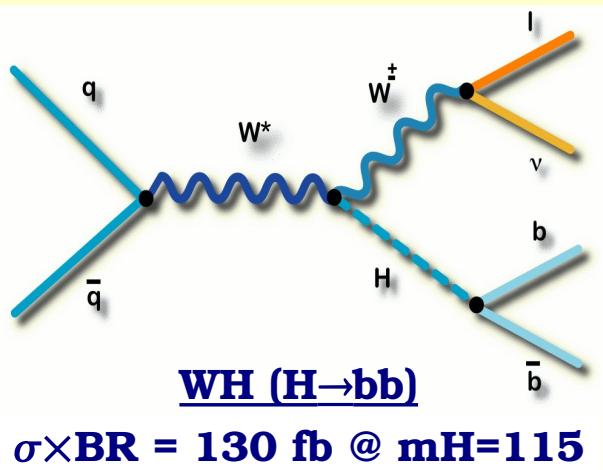
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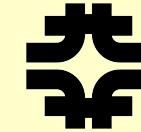
Backgrounds & Rates



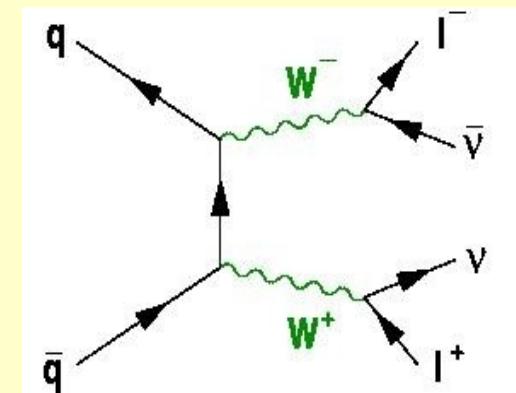
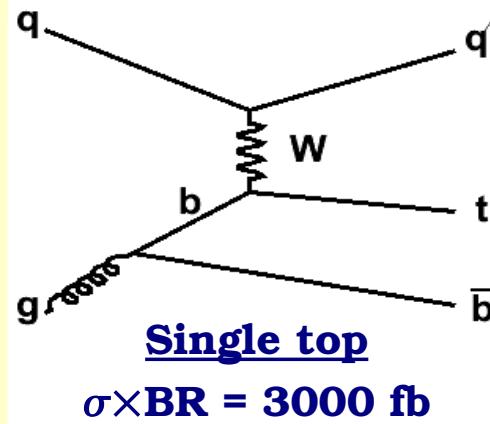
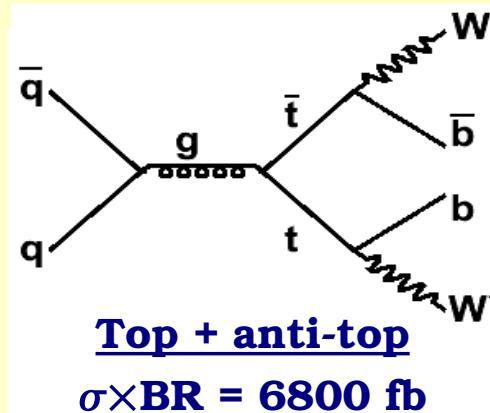
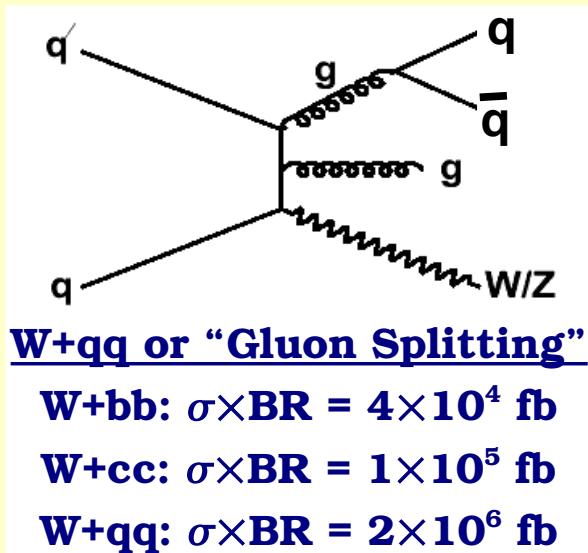
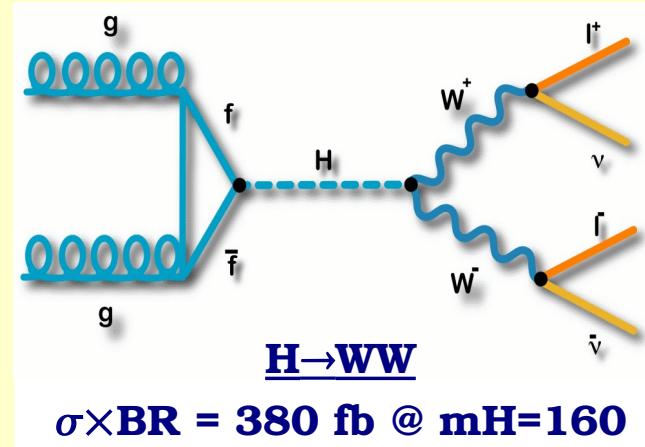
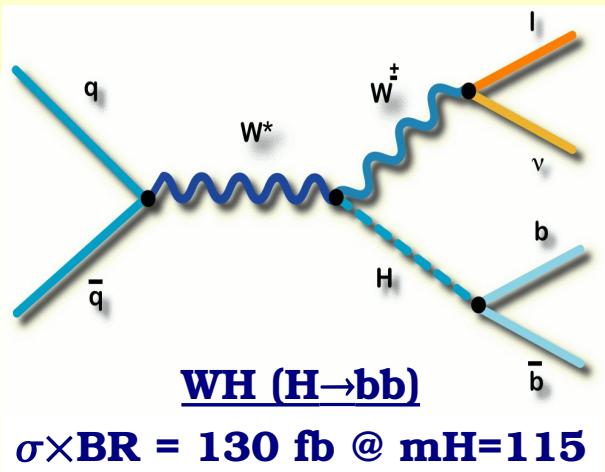
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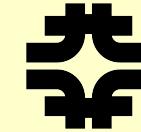
Backgrounds & Rates



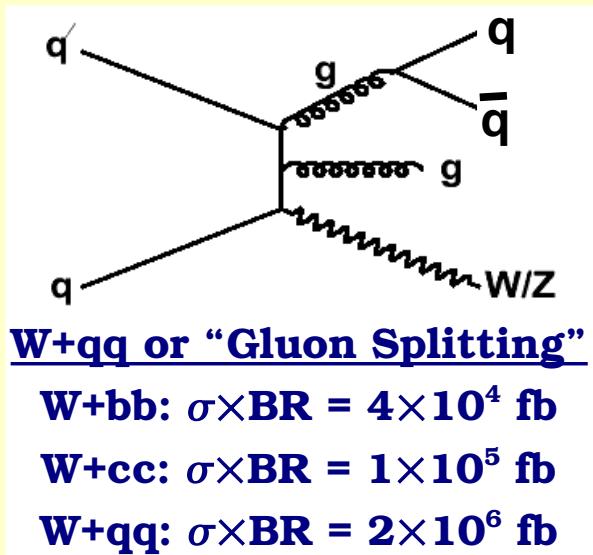
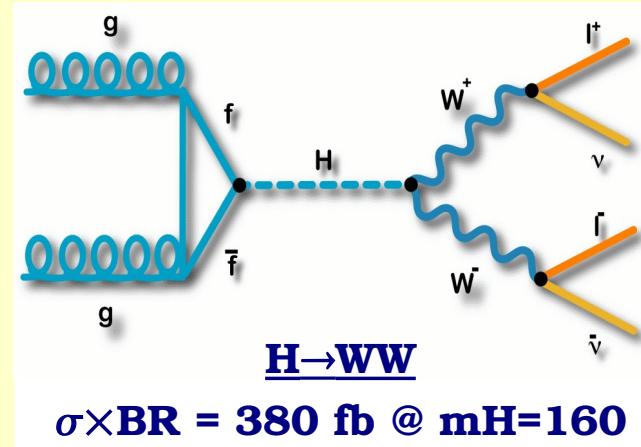
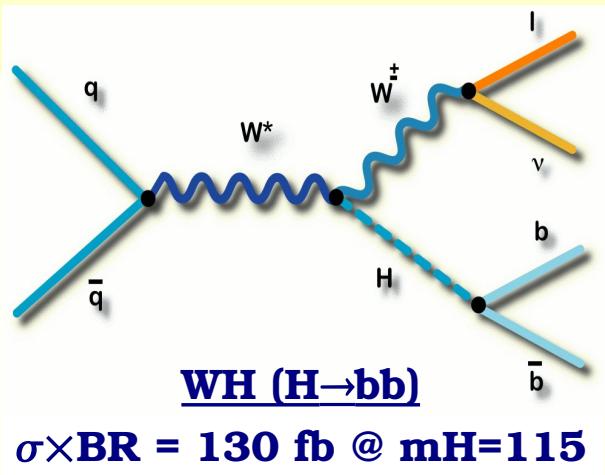
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Backgrounds & Rates



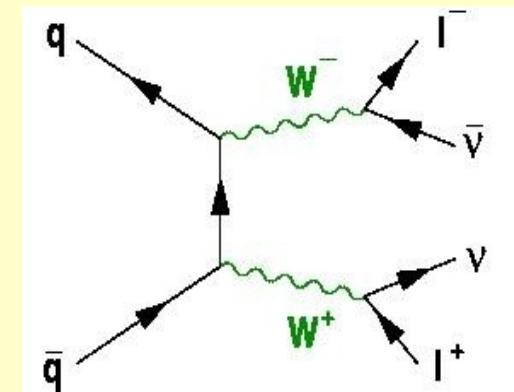
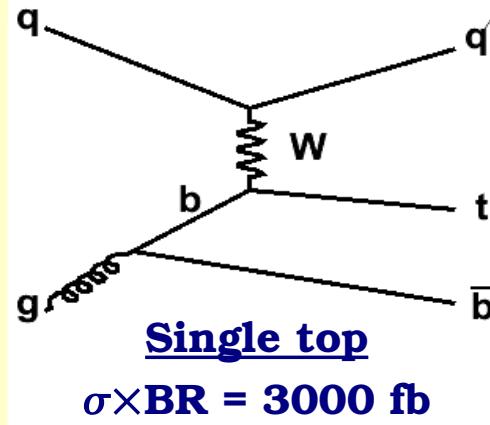
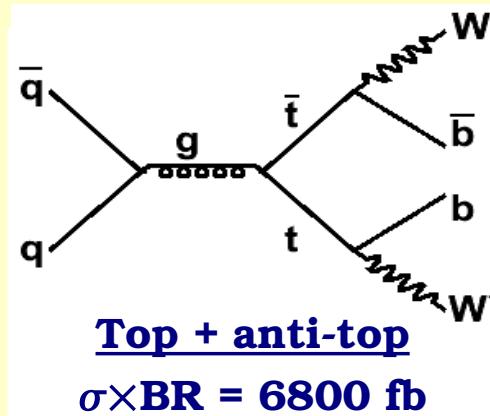
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**Large rate of QCD multi-jet production
("fake" lepton + neutrino)**

Simulations notoriously poor

⇒ **Measured in data**



WW: $\sigma \times BR = 13000 \text{ fb}$

WZ: $\sigma \times BR = 4000 \text{ fb}$

ZZ: $\sigma \times BR = 1500 \text{ fb}$

Searching for $H \rightarrow bb$

- Rely on $W \rightarrow l\nu$ and $Z \rightarrow ll$ provide handle to reduce backgrounds

Start with high-pT leptons: $p_T > 20$ (**15**) GeV

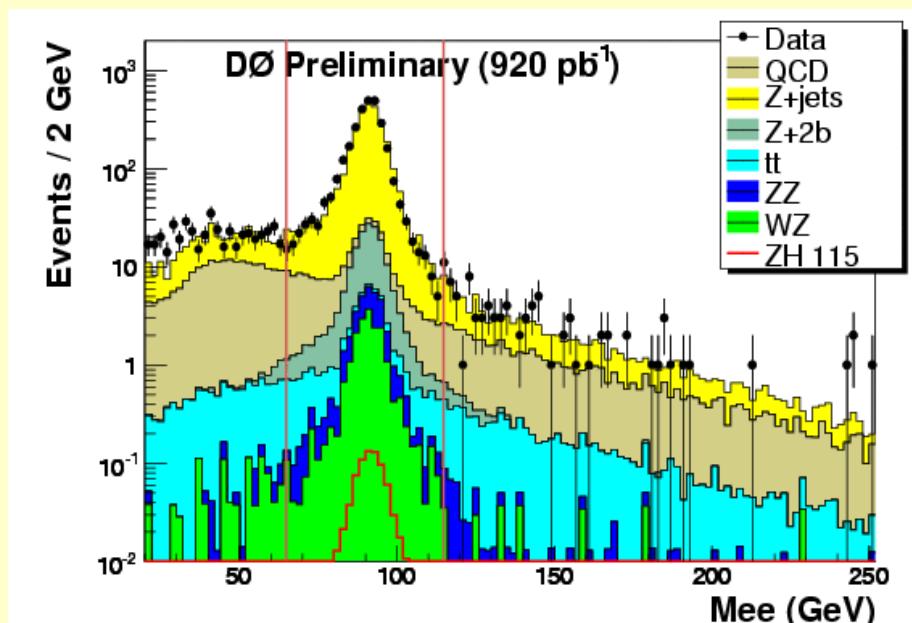
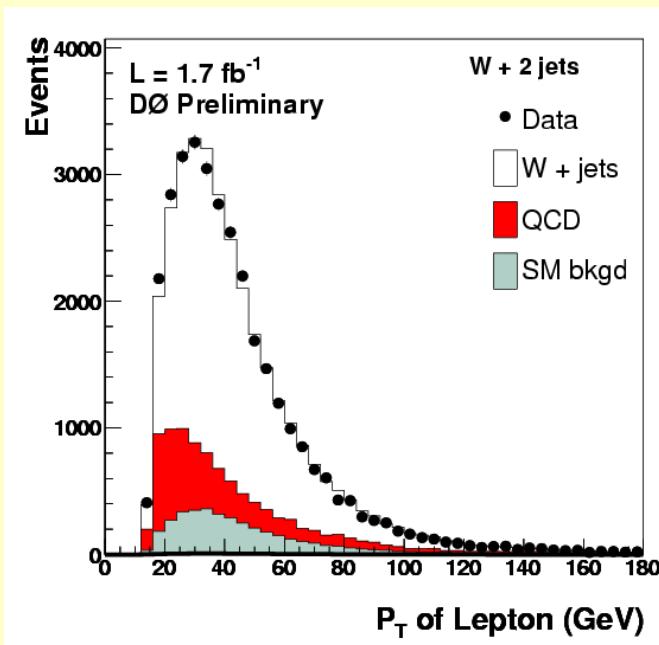
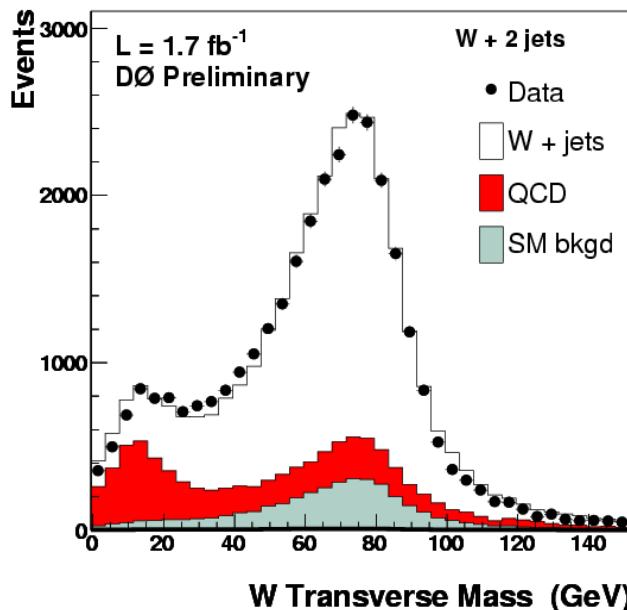
Neutrinos manifest as missing transverse energy

$WH \rightarrow l\nu bb$: MET > 20 GeV

$ZH \rightarrow ll bb$: MET *should be small!!*

Vector boson resonant mass is very important

$$M_W^{trans} = \sqrt{2 p_T^l M E_T (1 - \cos(\Delta\phi))}$$



Searching for $H \rightarrow bb$

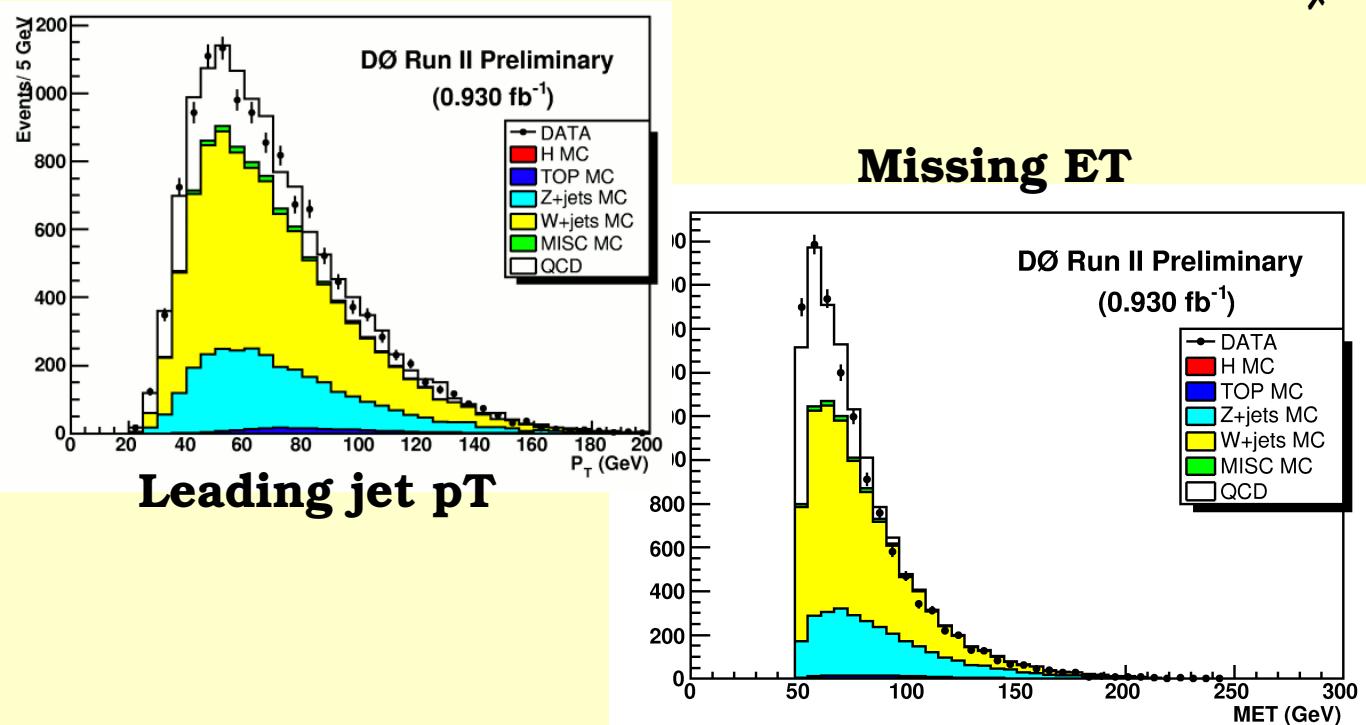
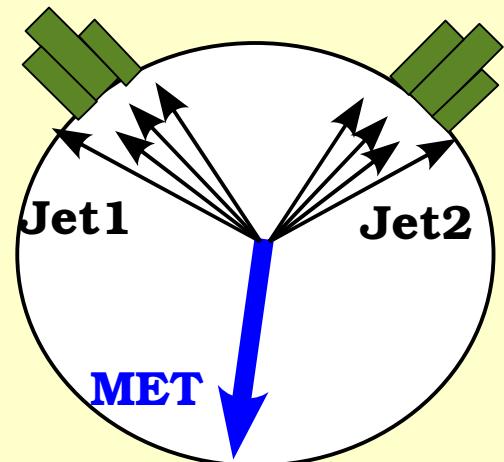
- For $ZH \rightarrow \nu\nu bb$ the search is more difficult:
no charged leptons!

Rely on large MET (neutrinos!)

“Physics” Bkgd: Z+jets, W+jets, top-pair, ZZ, WZ

“Instrumental” Bkgd: QCD multijets with
mismeasured jets

⇒ Asymmetry between MET and jet recoil isolates QCD



- Background reduction:

Trigger on large missing HT
(vector sum of jet ET),
select large MET:>50 GeV

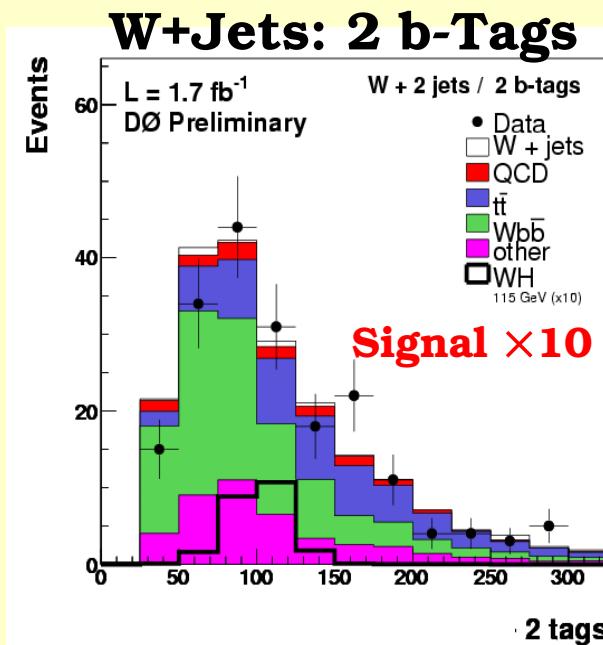
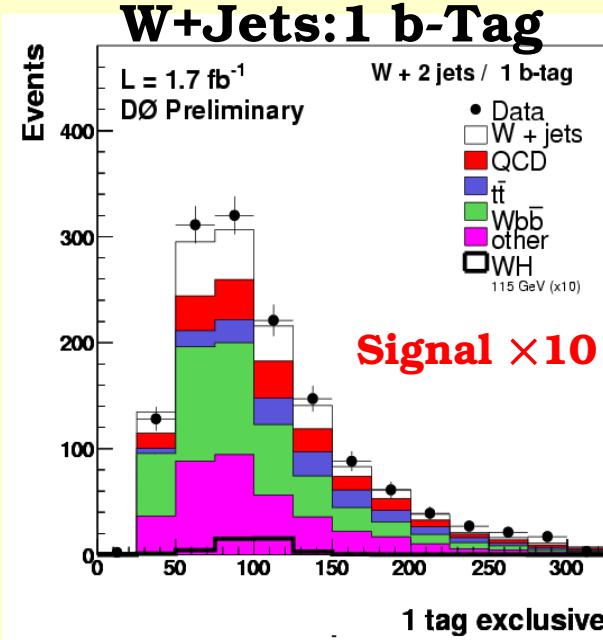
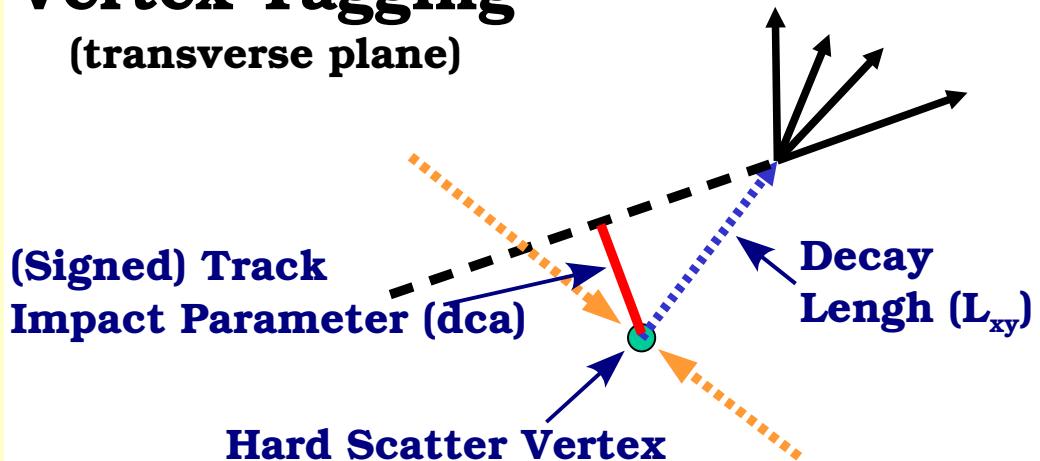
Select two high-pT jets to
define final state
($p_T > 20 \text{ GeV}$, $|\eta| < 2.5$)

Veto back-to-back jets: $\Delta\phi < 165^\circ$

Searching for $H \rightarrow b\bar{b}$

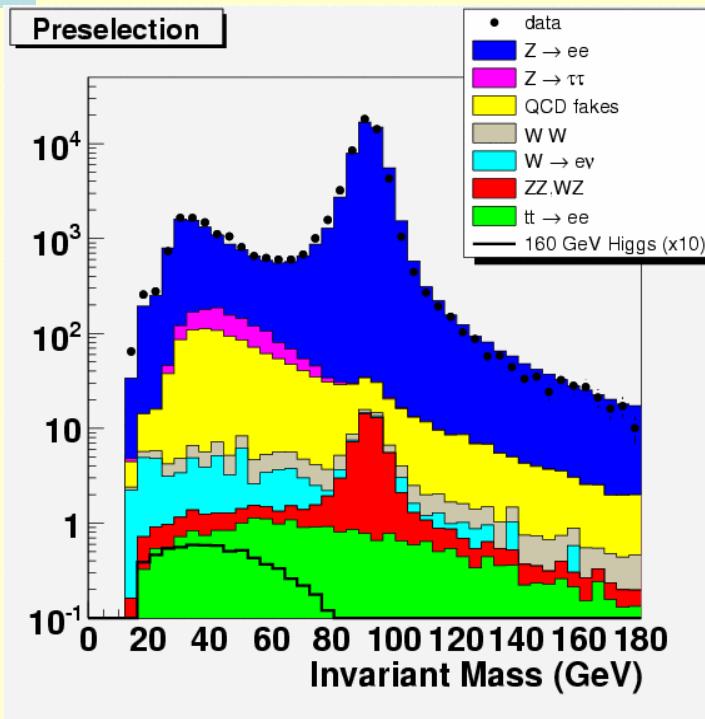
- ✗ Select quark-jets consistent with Higgs decays
 - At least two high-pT jets ($pT > 20-25$ GeV)
 - Angular separation consistent with scalar particle
 - Dijet invariant mass provides great discrimination
- ✗ Remove light-flavor jets via heavy-flavor “tagging”
 - Identify via displaced vertex due to b-meson lifetime
- ✗ Two main categories
 - Impact parameter based
 - Secondary vertex based

Vertex Tagging (transverse plane)

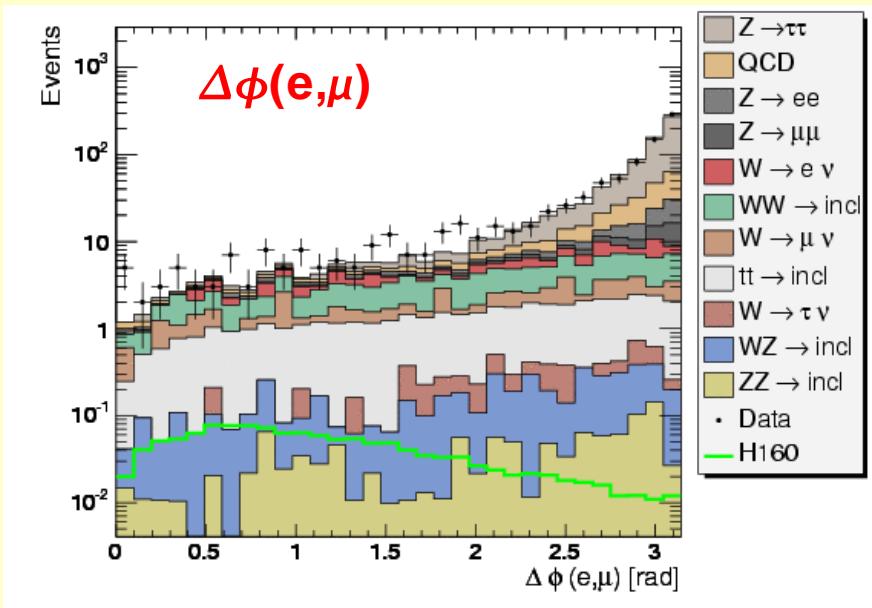
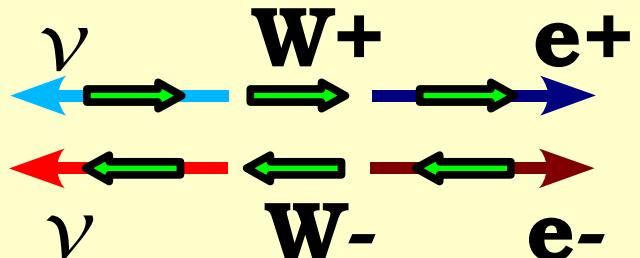


Searching for $H \rightarrow W^+W^-$

- x Select only leptonic W-boson decays
 - Select high-pT leptons (electrons & muons)
 - Use Z-peak for normalization ($ee/\mu\mu$), veto region near peak
 - Require large missing transverse energy signature from neutrinos, restrict sum of MET + lepton pT (scalar and vector)
- Preselection removes W/Z+jets, but Standard Model W^+W^- production remains

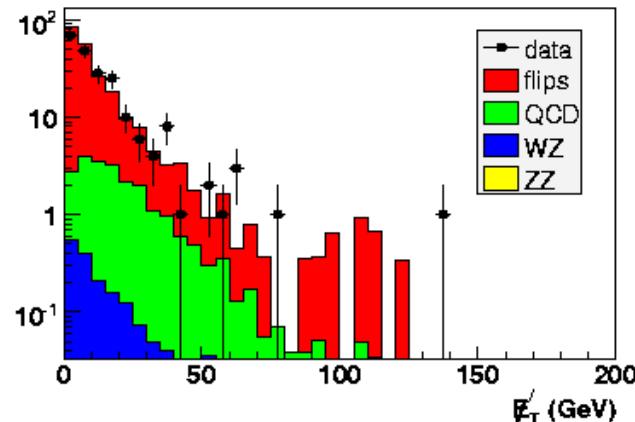
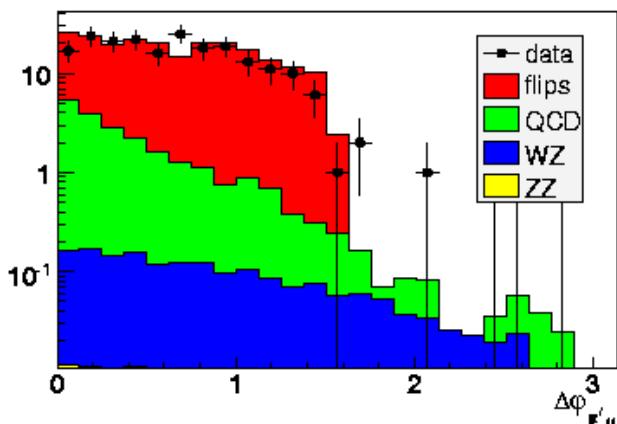
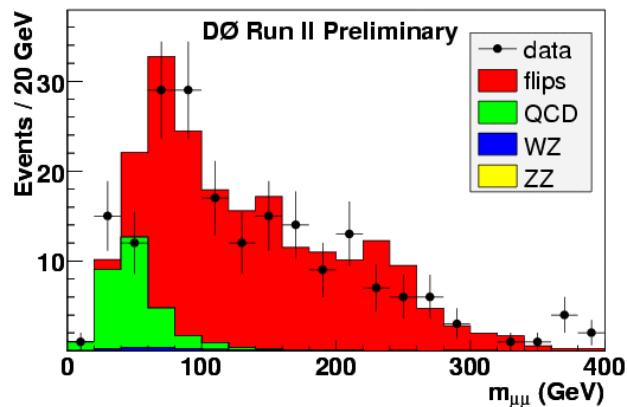


- x Scalar higgs (spin-0) provides natural discrimination due to spin correlation
- Leptons prefer to be collinear



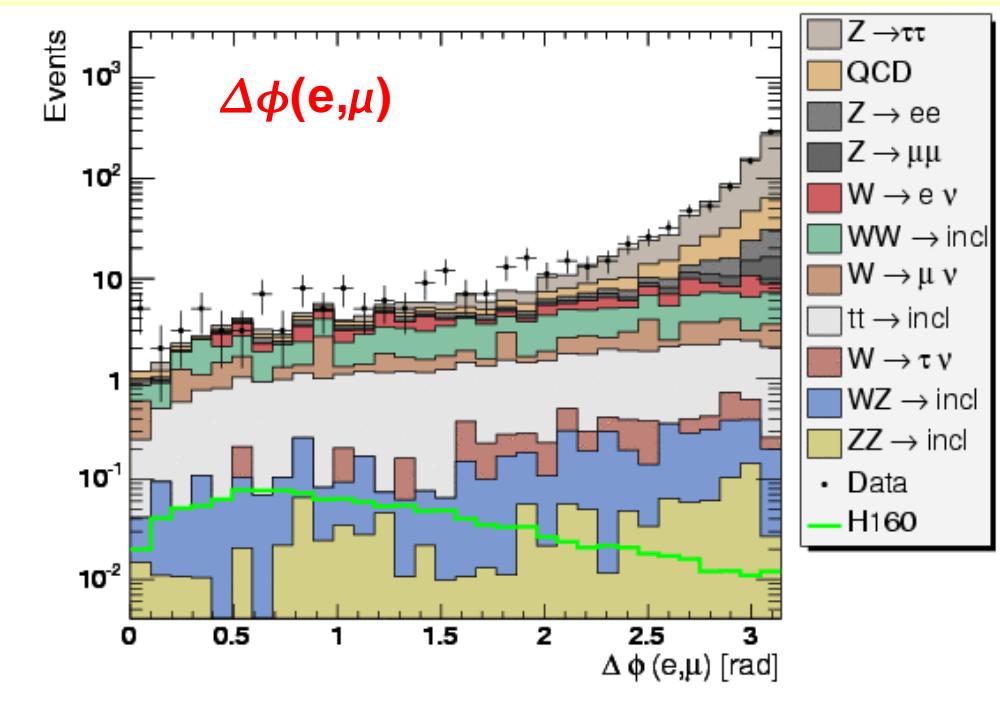
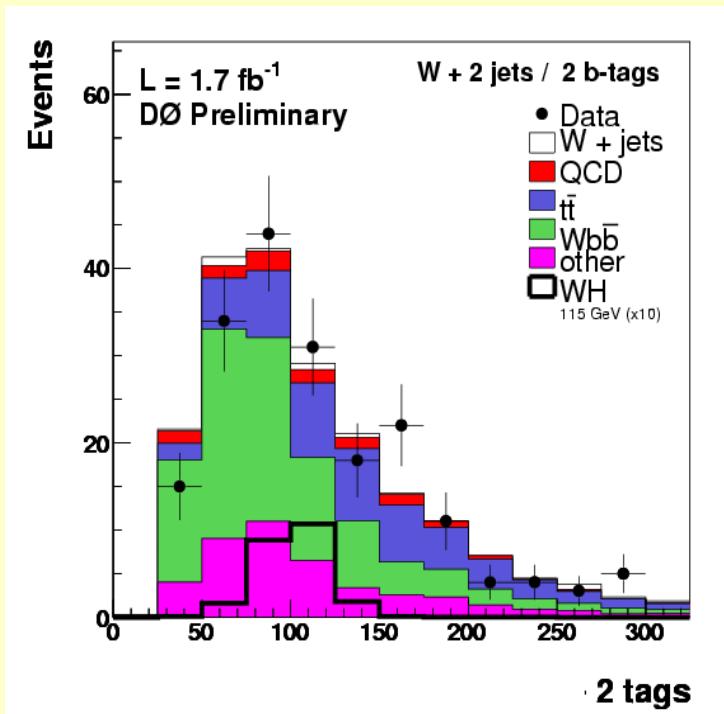
Searching for $WH \rightarrow WW^+W^-$

- ✗ Very rich spectrum in leptons
 - Three W-bosons give 1-, 2-, 3-lepton final states
 - First version of analysis reduces backgrounds by selecting same-signed lepton final states
 - Di-lepton mass, angles, MET all good discriminators
- ✗ $WH \rightarrow l^\pm l^\pm + X$ has few backgrounds
 - Physics Bkgds: WZ, ZZ, W+jets (*fake lepton*)
 - Instrumental Bkgds: QCD (*fake leptons*), charge-sign “flips”
- ✗ Instrumental backgrounds carefully measured
 - Flips: Use multiple sign measurements for orthogonal “non-Z” events. Diagonalize matrix of Same-Same, Opposite-Same, & Opposite-Opposite events



Final Variables

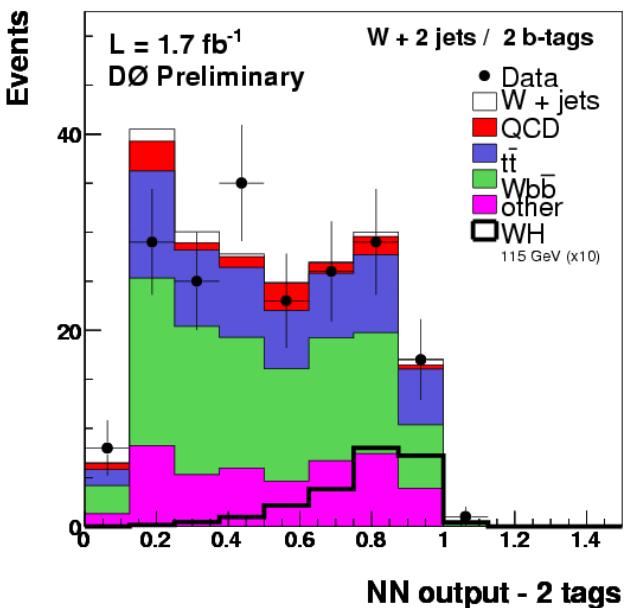
- ✗ In the past, analysis challenges forced analyses to estimate search sensitivity from single kinematic variables
 - Data/MC agreement, sample statistics, background uncertainties
- ✗ Significant reduction in potential search power
 - Low mass: Dijet invariant mass
 - High mass: Dilepton angular separation



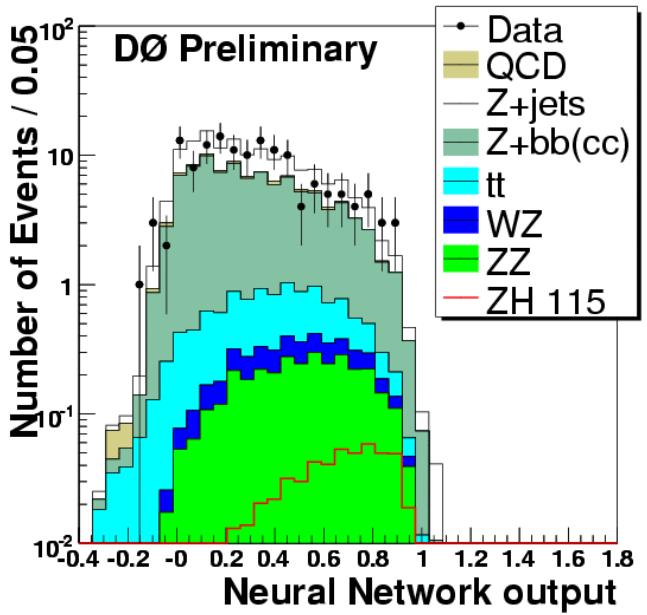
Advanced Analysis Techniques

- × Several multivariate approaches have emerged
- Neural networks, likelihood discriminants, Matrix Element discriminants
- Techniques demonstrate 20-50% improvement in signal isolation
- Analysis refinement & optimization still ongoing*

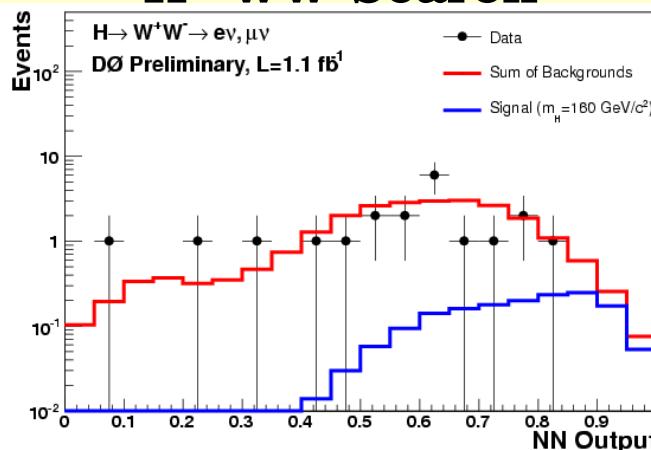
WH Search



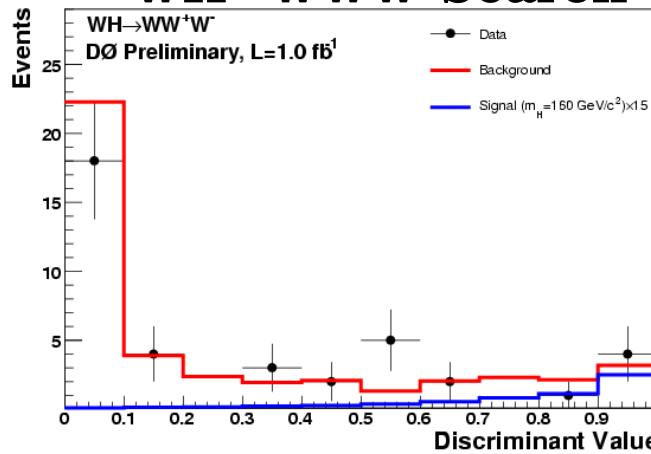
ZH Search



H \rightarrow WW Search



WH \rightarrow WWW Search



Advanced Analysis Techniques

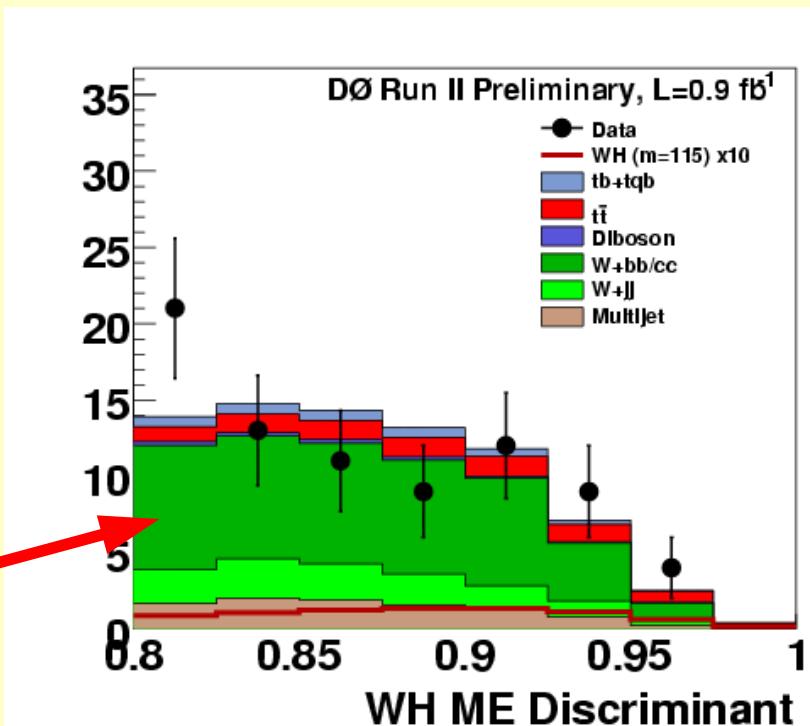
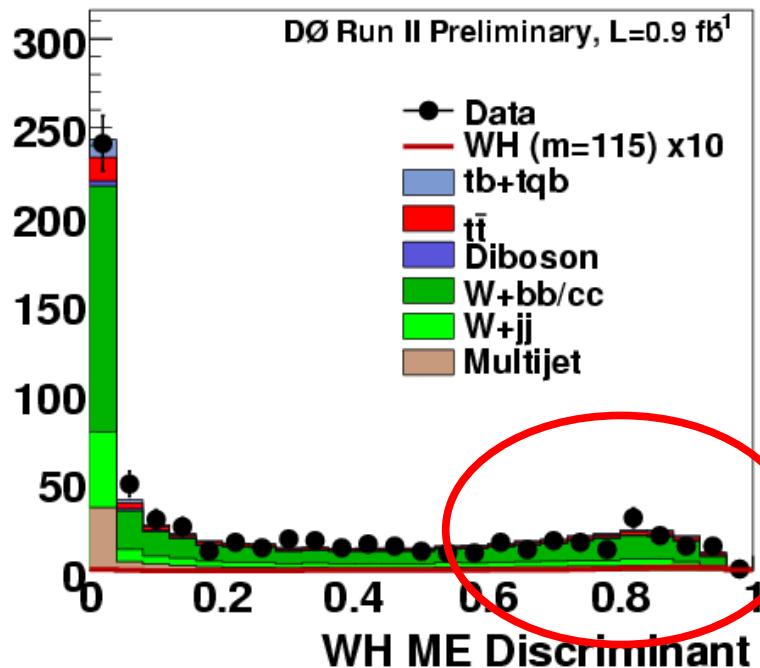
- “Matrix Element” analysis technique

Developed in Tevatron top-quark search/measurements

Utilized in both CDF, DZero single top quark & Higgs analyses as well

Use signal/bkgd production Matrix Elements (tree-level) to form likelihood discriminant: *over 35% improvement in sensitivity!*

Can be combined with complementary analysis techniques



Summary of Analyses

- ✗ WH \rightarrow lvbb: ev, mv, single-tag, double-tag
 - RunIIA+RunIIB: 1.7 fb^{-1} (4 analyses)
- ✗ ZH \rightarrow vvbb: ZH and WH(missing lepton) signals
 - RunIIA: 0.95 fb^{-1} (1 analysis)
- ✗ ZH \rightarrow llbb: ee + mm, single-tag, double-tag
 - RunIIA: 1.1 fb^{-1} (4 analyses)
- ✗ H \rightarrow WW: ee, em, mm
 - RunIIA+RunIIB: 1.7 fb^{-1} (3 analyses)
- ✗ WH \rightarrow WWW: ee, em, mm
 - RunIIA: 1.0 fb^{-1} (3 analyses)

15 analysis channels, 11 Higgs decays!

Testing the Higgs Hypothesis

- ✗ Assuming we've assembled all the pieces properly, proceed to test our two hypotheses: **H0** and **H1**
- ✗ Tevatron \Rightarrow utilize both Bayesian and as-Frequentist-as-possible treatments
 - Uncertainty on nuisance parameters generally only has a Bayesian interpretation
 - Access to the subsidiary measurement isn't guaranteed, generally nuisance parameters are some number for which we apply a non-Frequentist belief distribution
- ✗ Bayesian treatment: choose an appropriate prior and integrate likelihoods over nuisance parameters
 - Choice of prior: flat in non-negative number of expected signal events (Heaviside step function / Jeffrey's Prior)
- ✗ Semi-Frequentist approach (CLs): use likelihood ratio as a test statistic and simulate pseudoexperiments to populate distributions for **H0** and **H1**
 - Vary nuisance parameters according to their Bayesian credibility distributions in the pseudoexperiments to obtain the “prior predictive ensemble”

We find CLs usually gives slightly better sensitivity, and fewer coverage issues



The CLs Approach

- ✗ Adopt Bayes-type definition of confidence levels

This normalization can help remove ambiguity of poorly modeled bkgds

$$CL_s = 1 - \frac{CL_{s+b}}{CL_b}$$

Not a discovery tool!

- ✗ Use likelihood ratio for search test statistic

Distributions of simulated outcomes are populated via Poisson trial with mean values given by b-only or s+b

Systematics are sampled for each Poisson trial, correlations propagated

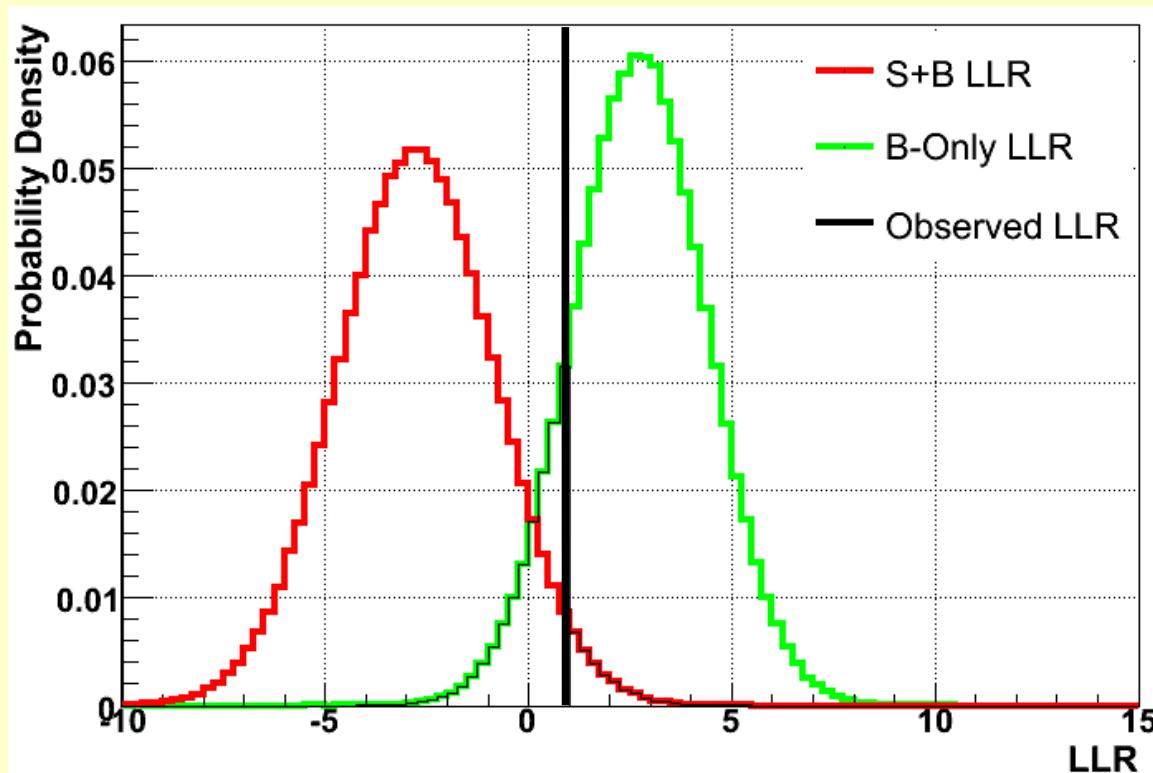
$$Q(\vec{s}, \vec{b}, \vec{n}_{cand}) = \prod_{i=0}^{N_{Channels}} \prod_{j=0}^{N_{bins}} \frac{(s+b)_{ij}^{n_{ij}} e^{(s+b)_{ij}}}{n_{ij}!} / \frac{b_{ij}^{n_{ij}} e^{b_{ij}}}{n_{ij}!}$$

Log likelihood ratio monotonic in number of candidate events

$$-2\ln(Q) = 2s_{tot} - 2 \sum_{i=0}^{N_{Channels}} \sum_{j=0}^{N_{bins}} n_{ij} \ln \left(1 + \frac{s_{ij}}{b_{ij}} \right)$$

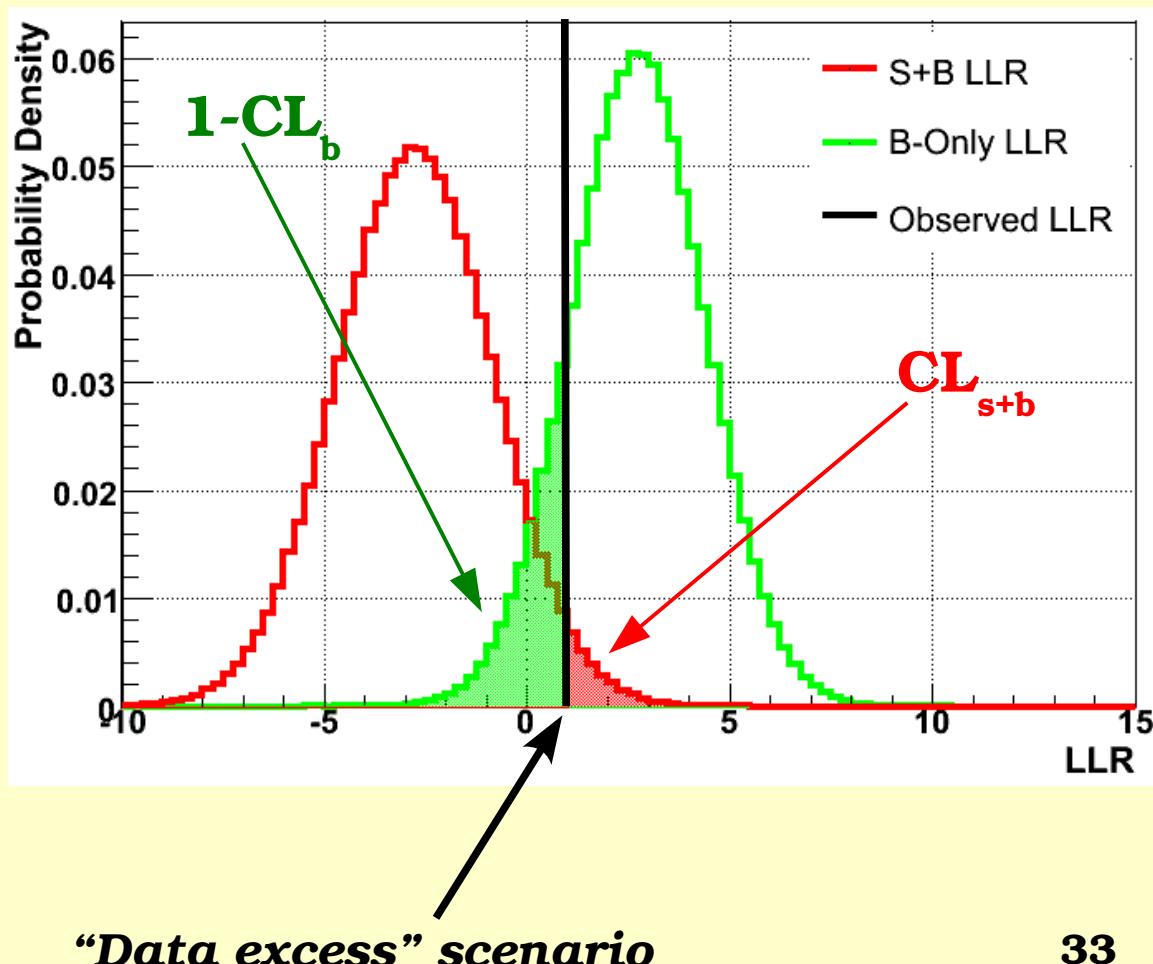
- x Model repeated outcomes of the experiment via Poisson distribution
 - Simulate Signal+Bkgd and Bkgd-Only outcomes based on predictions
 - Uncertainties on nuisance parameters folded in via Gaussian smearing

- x **Black line: Observed LLR value (LLR_{obs})**
Determined by data
- x **Green: Bkgd-only hypothesis**
 CL_b is region to right of LLR_{obs}
Equals ~50% for good bkgd/data agreement
- x **Red: Signal+bkgd hypothesis**
 CL_{s+b} is region to right of LLR_{obs}



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Profile Likelihood

- Now we can get serious about systematics and limits

First define the best fit of our MC model to data

Assume background transforms as:

$$\hat{X}_i = X_i \prod_k (1 + \sigma_i^k S_k)$$

Where S_k has a mean of 0 and width of 1 (*unit Gaussian*)

Our “*educated guesses*” to size of uncertainties must be taken into consideration along with data/MC shapes

These can (should) be limited via our previous discussion

Need to change our estimator to penalize departures from input distributions

Minimize estimator by varying S_k values!

Likelihood now a function of nuisance parameters

$$\chi^2 = 2 \sum_i (\hat{X}_i - D_i)^2 - D_i \ln \left(\frac{\hat{X}_i}{D_i} \right) + \sum_k S_k^2$$

Sidebands in CLs

- ✗ The CLs approach does not have sideband constraints *a priori*
 - ✗ But we'd still like to benefit, otherwise significance can be severely degraded via marginalization
 - ✗ Several solutions exist in statistical literature, we choose to maximize likelihoods over nuisance parameters:

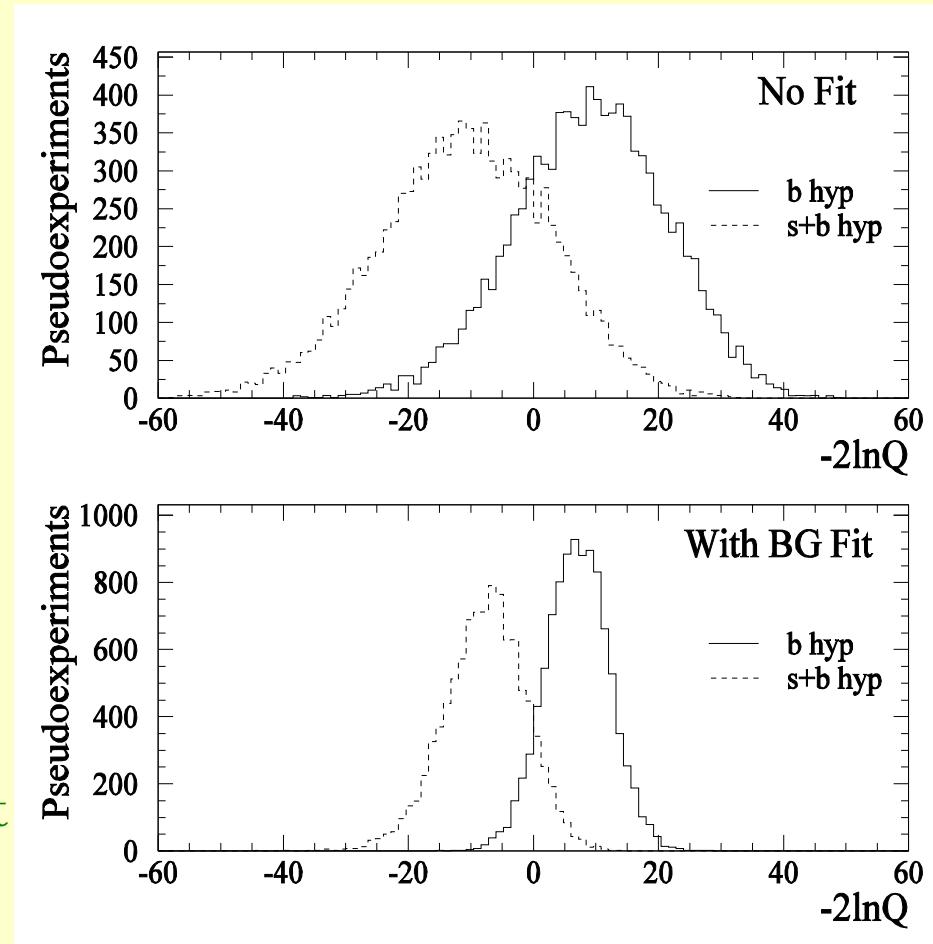
$$Q = \frac{L(x|\theta_{R1}, \hat{\theta}_S)}{L(x|\theta_{R0}, \hat{\hat{\theta}}_S)}$$

Two maximizations performed independent of physics parameters, once for each pseudoexperiment

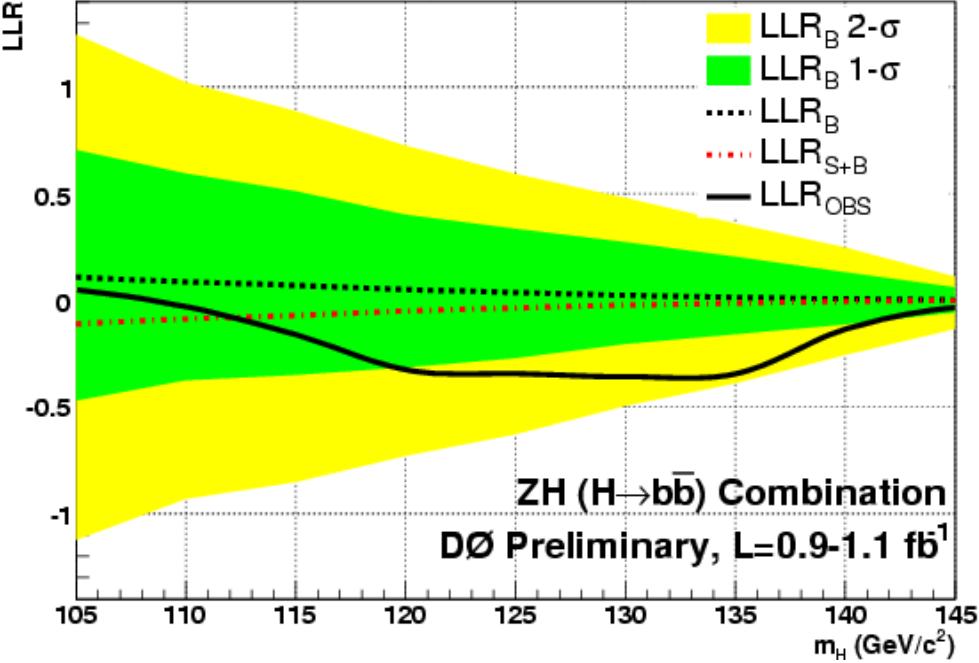
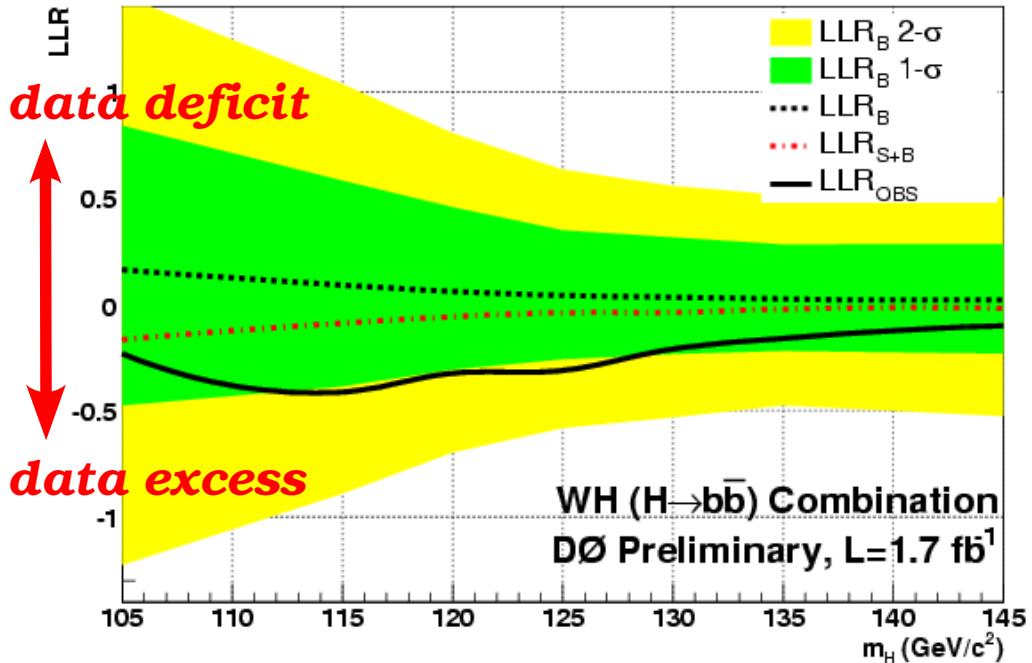
- ✗ θ_{R1} : Physics parameters in **H1**, θ_{R0} : Physics parameters in **HO**
- ✗ $\hat{\theta}_S$: Nuisance parameters which maximize L for **H1**
- ✗ $\hat{\hat{\theta}}_S$: Nuisance parameters which maximize L for **HO**
 - ✗ If sidebands are unambiguous, optionally maximize just once per pseudoexperiment – introduces bias in sideband selection

Maximum Likelihood Fits

- ✗ Nominally, finding best fit for nuisance parameters makes both **H1** and **H0** match the data better
- ✗ In practice, improves separation of **H1** and **H0** in $-2\ln Q$ distributions: effective constraint on size of nuisance parameter uncertainty
 - ✗ We rely on MINUIT for fitting, assume same answer for identical problems
- ✗ Question: Doesn't this mean your uncertainties are overestimated?
- ✗ Answer: In many cases this is true, but subsidiary experiments for constraint are often unavailable.



Individual Channel Results

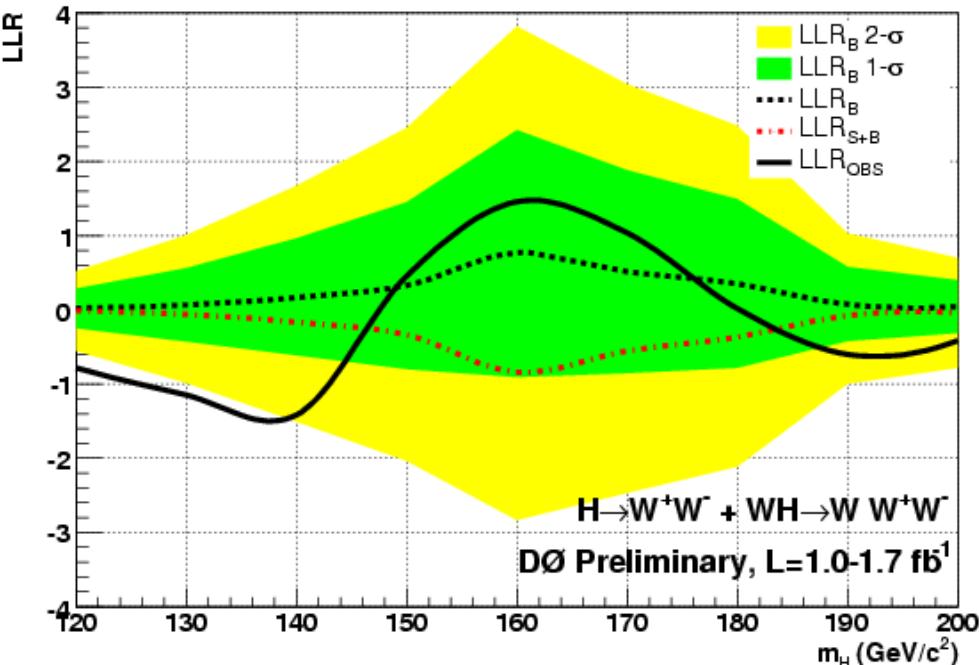


✗ LLR profile vs Higgs mass

Dashed lines show **S+B** and **B-Only**
mean value

Shaded bands indicate 1- and 2- σ
variation of B-only distribution

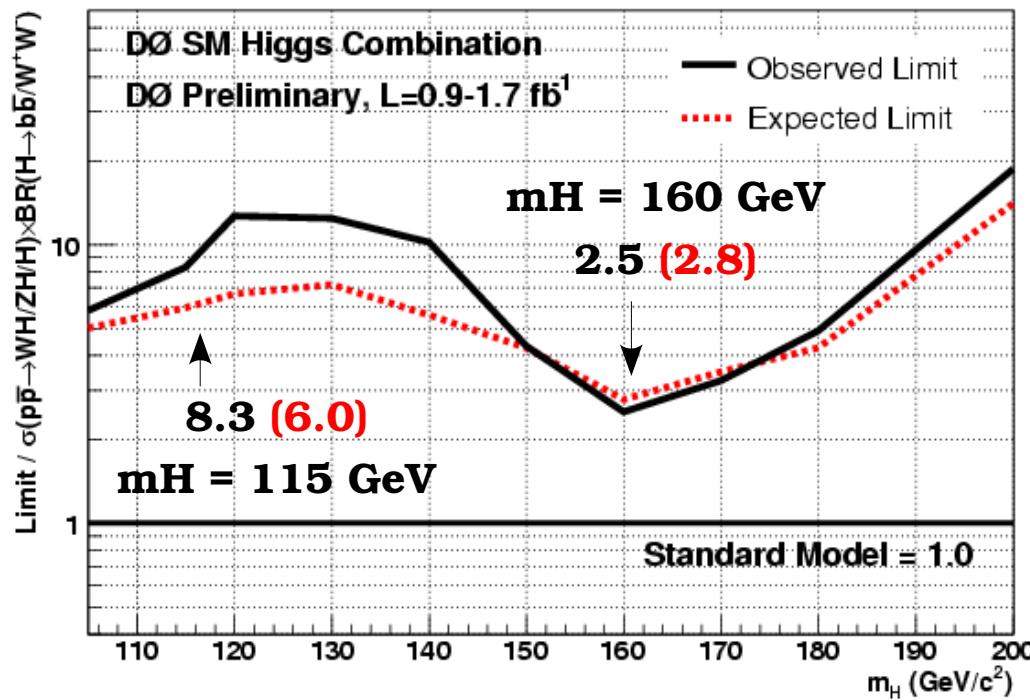
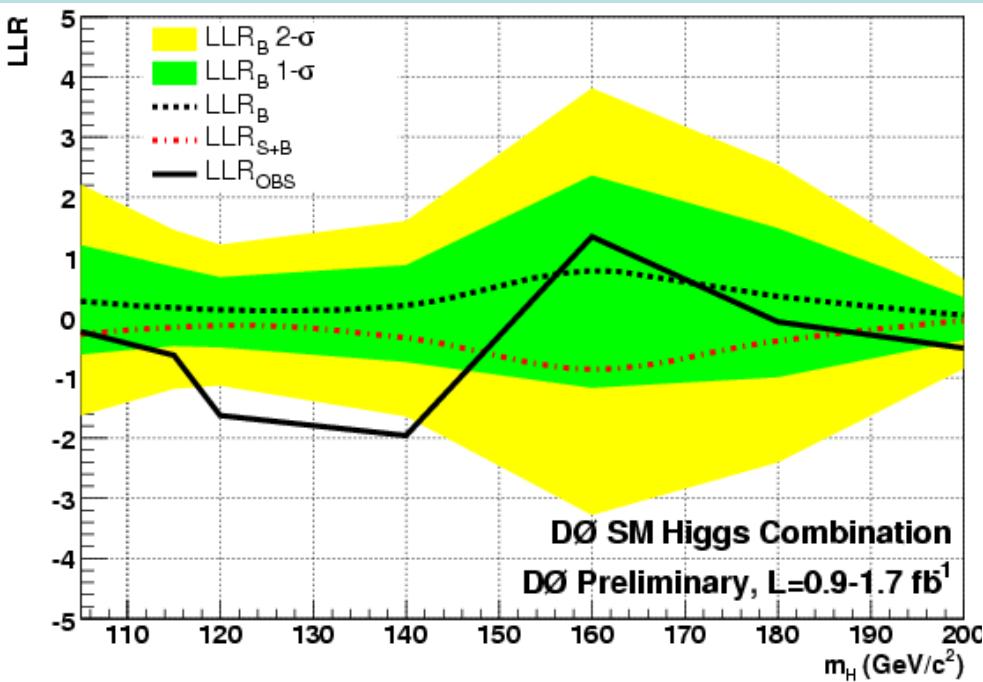
Solid black line indicates **data**
observation



Limits on Higgs Production



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Combined results for 15 analysis channels covering $0.9-1.7 \text{ fb}^{-1}$

Results reported as ratio of 95% CL limit on $\sigma \times BR$ to standard model expectation \Rightarrow Sensitivity to SM achieved when $R \leq 1.0$

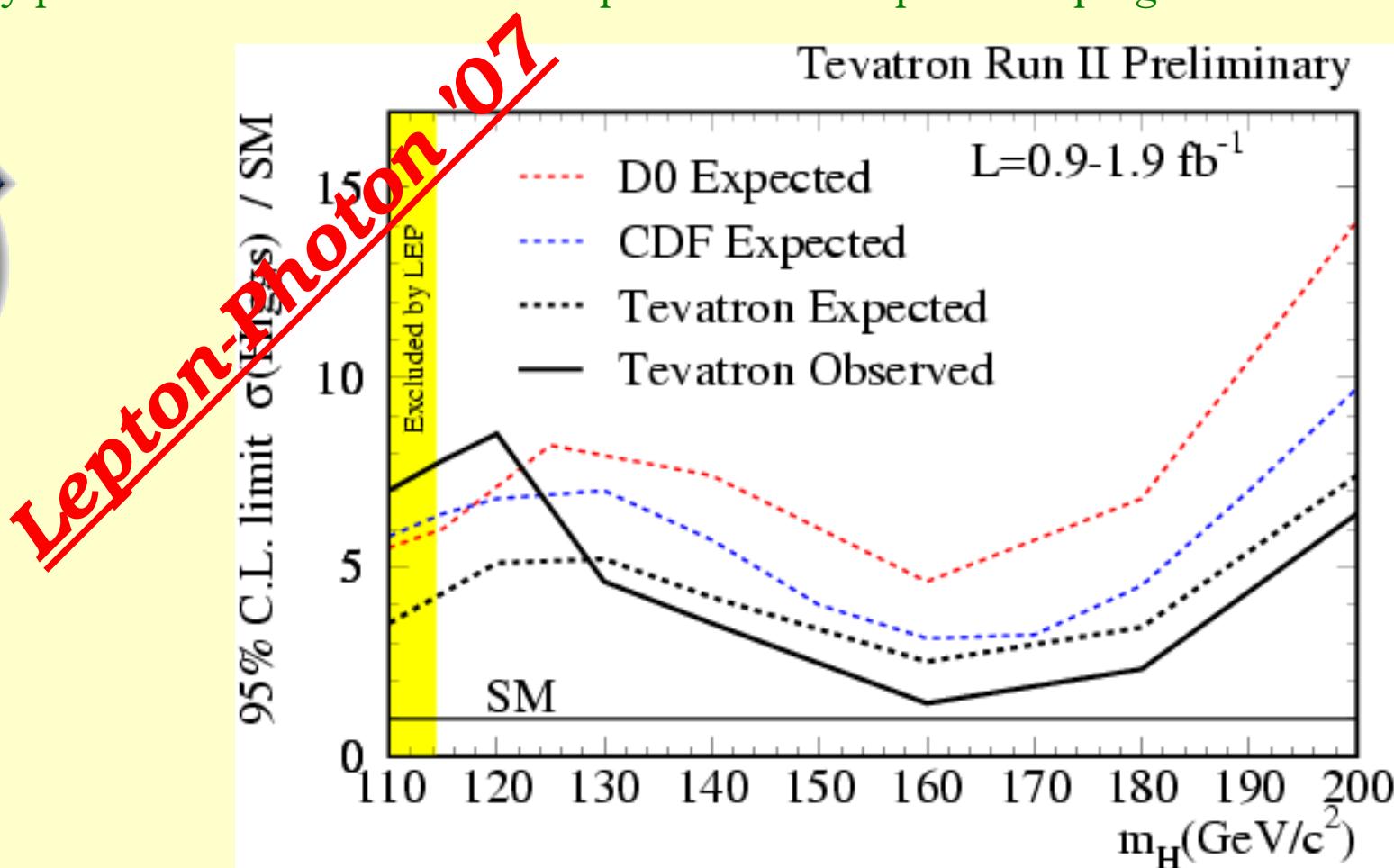
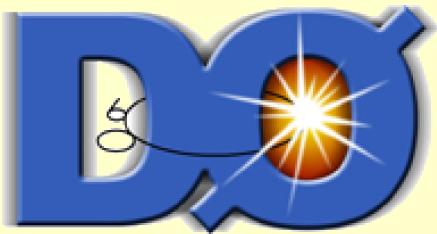
Analyzed luminosity depends on analysis channel & understanding of bkgds

Combined Tevatron Limits

CDF+DZero combination evaluated by the *Tevatron Higgs Working Group*

CDF adds similar search sensitivity in all analysis channels

Total luminosity per channel varies between experiments: “snapshot” of progress



Obs (exp) = 7.8 (4.3) @ $m_H = 115 \text{ GeV}$

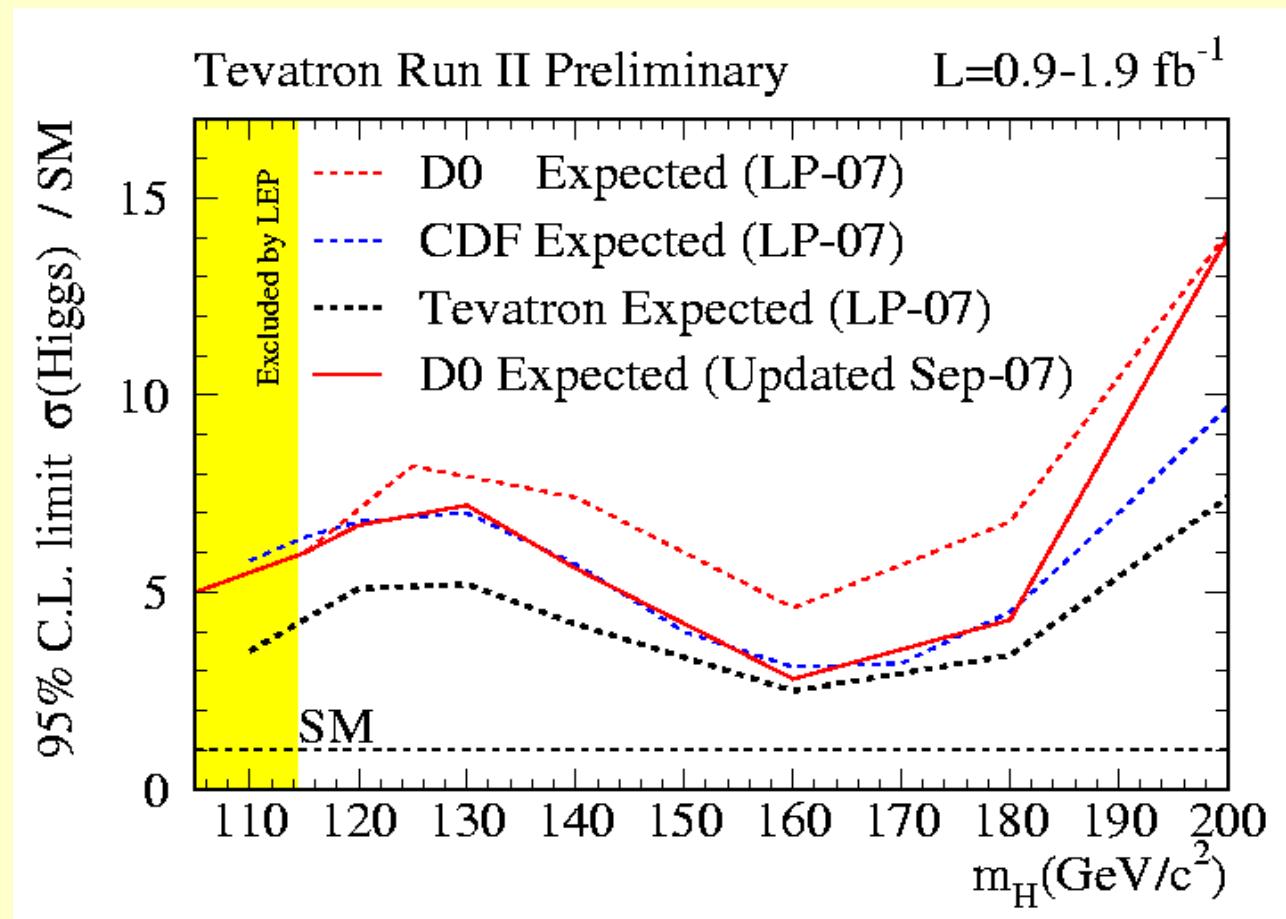
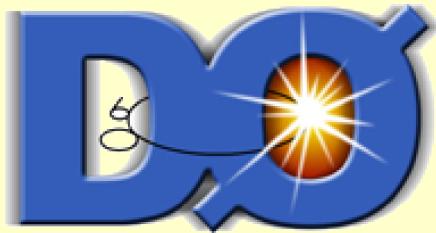
Obs (exp) = 1.4 (2.5) @ $m_H = 160 \text{ GeV}$

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Missing Pieces

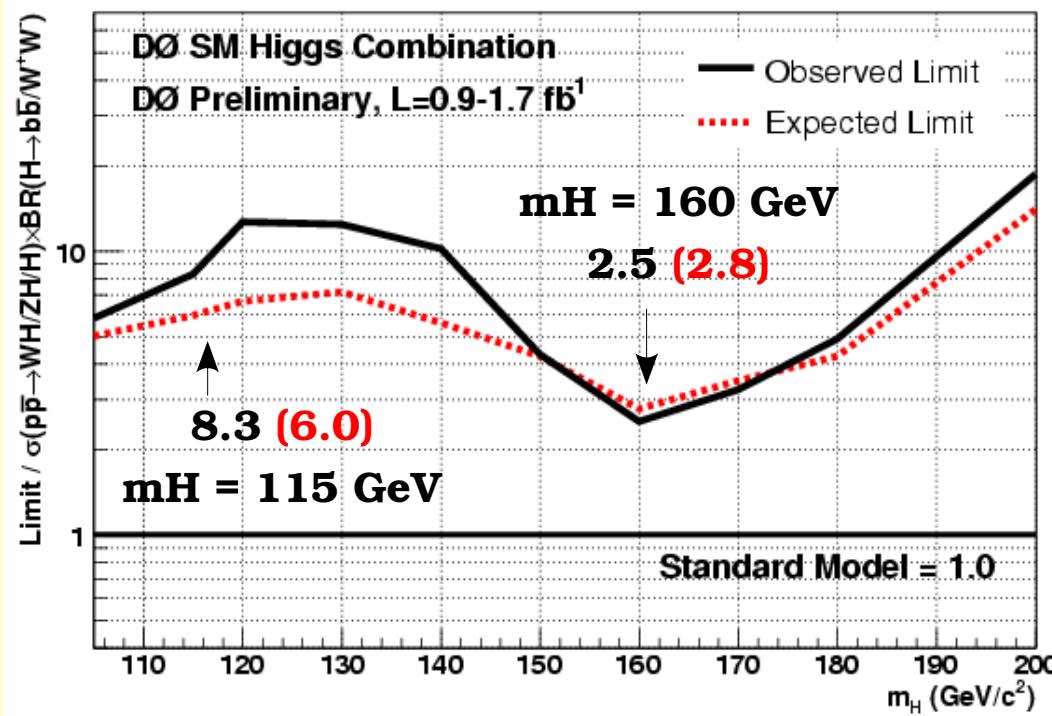
General Issues

Many channels systematics-limited:
growing dataset reduces uncertainties

More powerful analysis techniques
become feasible with growing dataset

Currently have only ~55% of desired
branching fraction

Recovery of efficiency via fiducial volume



Low Mass Higgs ($H \rightarrow bb$)

Dijet mass resolution drives sensitivity

ZH $\rightarrow \nu\nu bb$: Only using 0.9 fb⁻¹, updated
analysis in approval process

WH $\rightarrow \tau\nu bb$: Analysis under development

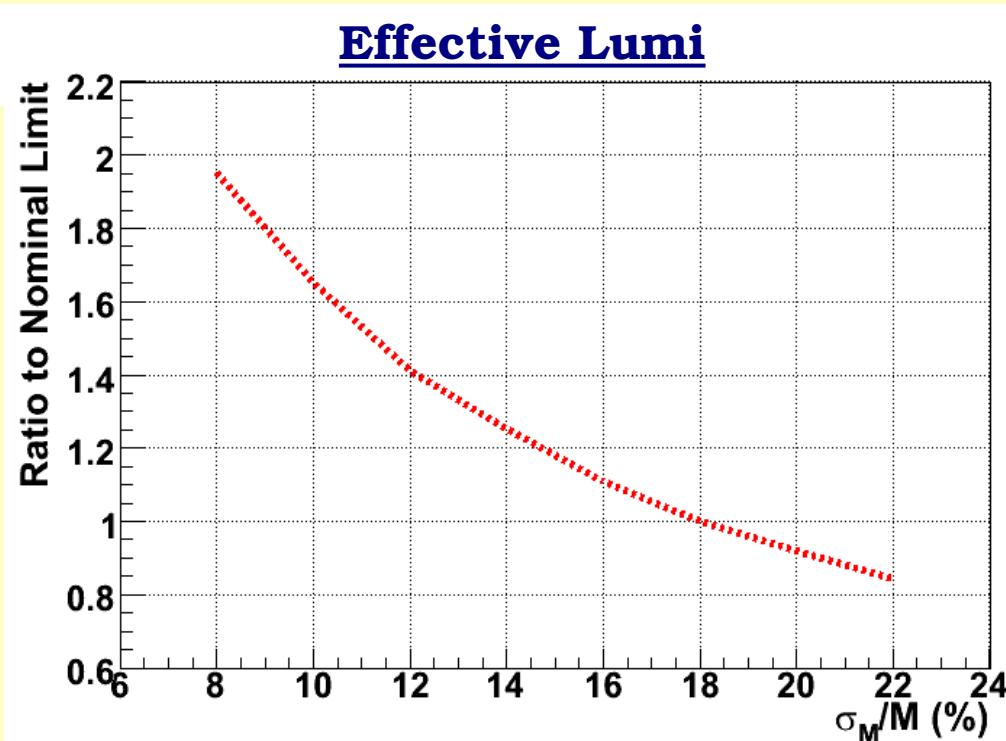
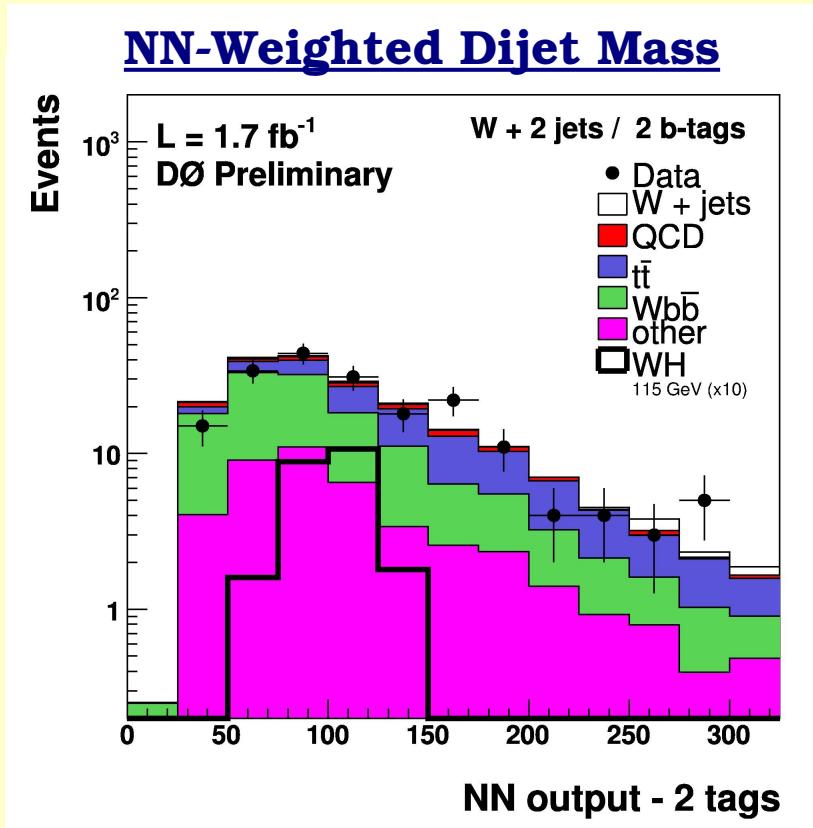
High Mass Higgs ($H \rightarrow WW$)

Initial neural network treatment
suboptimal

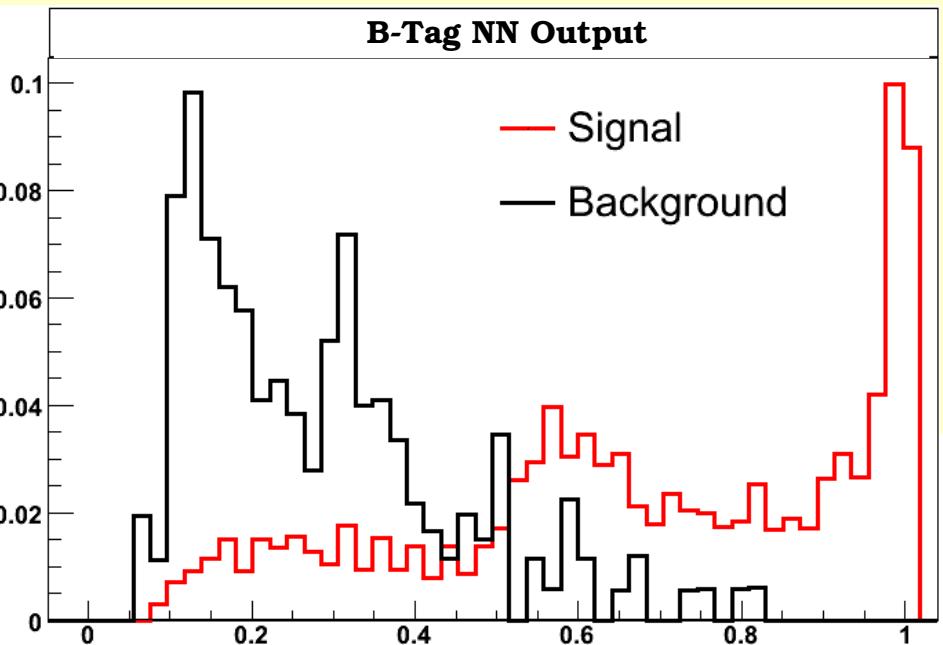
Tau channels and hadronic W decay
analyses under development

Major Challenges

- ✗ Low mass: dijet mass resolution
 - Still dominant variable in NN & ME analysis techniques
- ✗ Currently at ~18% \Rightarrow goal closer to 10%
 - Studies indicate ~14% is achievable soon
 - Largely determined by calibration of calorimeter*

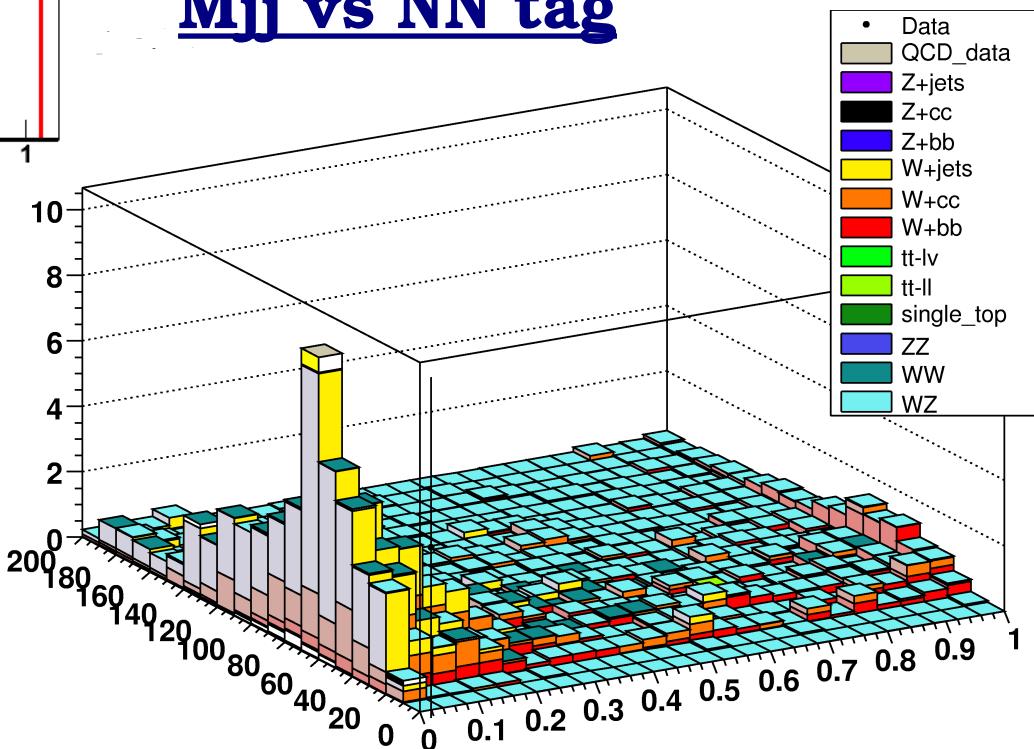


Major Challenges



✗ Low mass: b-Tagging

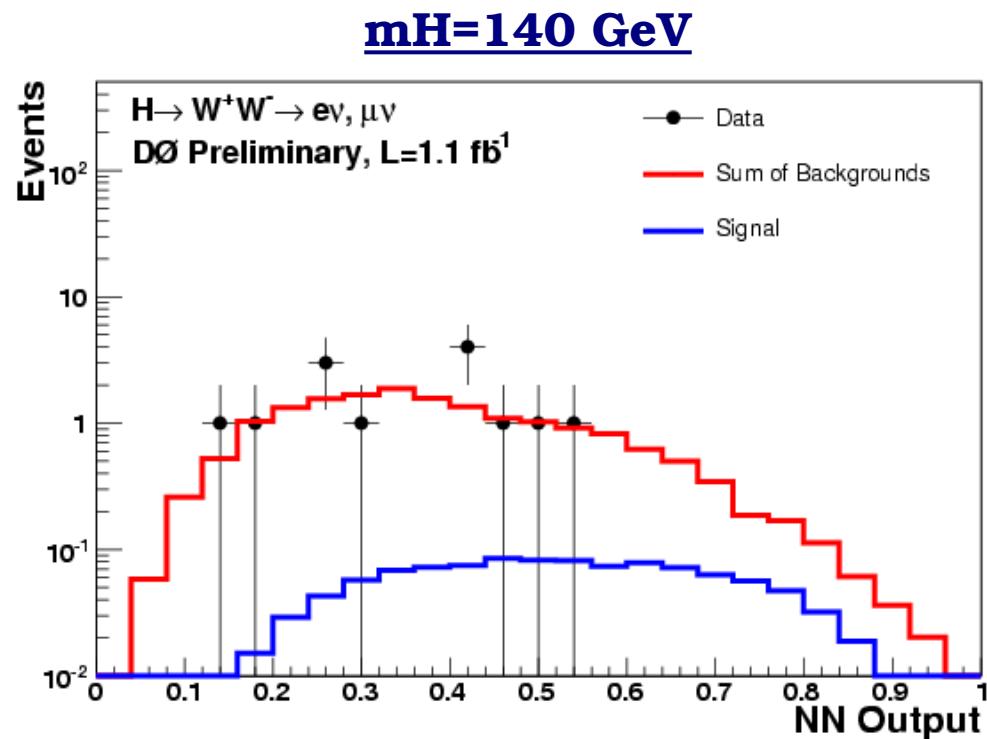
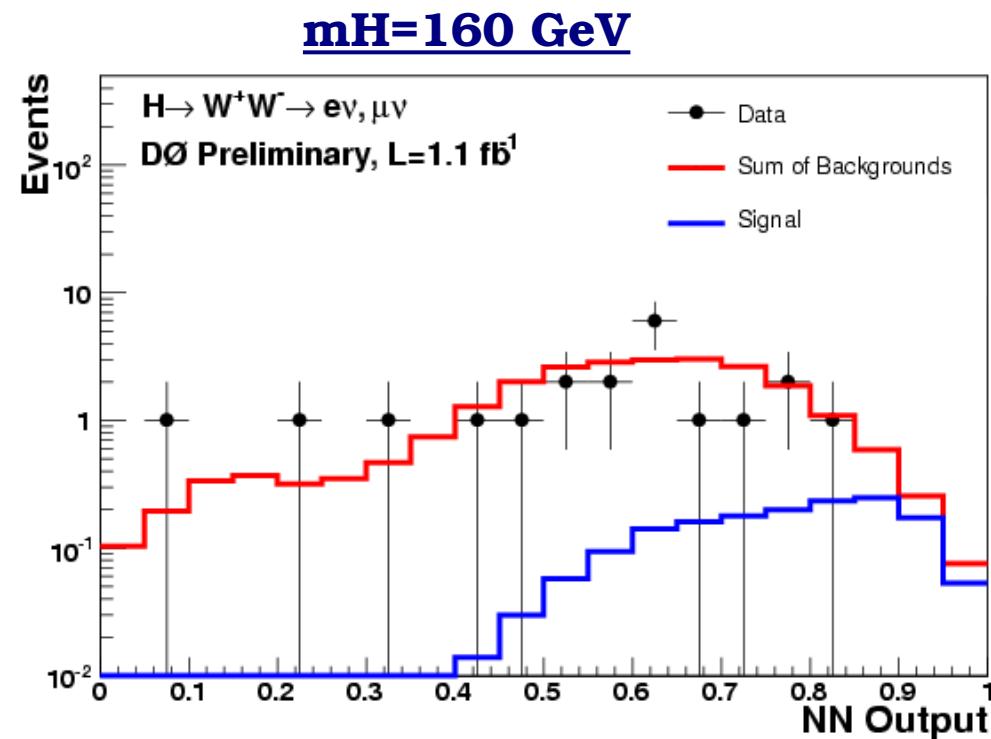
M_jj vs NN tag



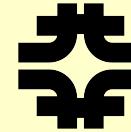
- ✗ Current b-Tag algorithm is NN based
 - NN shape not used in analyses due to modeling of tracking
 - Studies show ~30% improvement in sensitivity possible

Major Challenges

- ✗ High Mass: uniform background reduction
 - Signal kinematics change significantly over $100 < m_H < 200$ GeV
 - Dominant background(s) change as a function of mass
 - More data statistics needed for careful optimization
- ✗ Studies indicate 20% improvement possible for $100 < m_H < 140$ GeV



Sensitivity Projections



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Can project sensitivity as a function of analyzed luminosity

Analyzed luminosity = ~80% of delivered luminosity

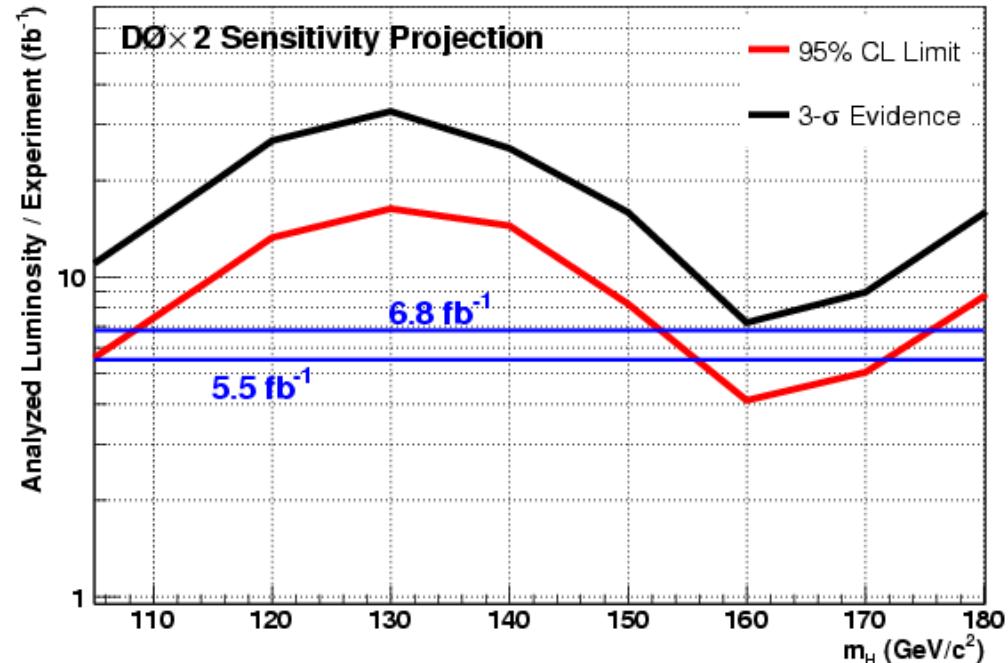
5.5 fb^{-1} = ~7 fb^{-1} delivered in 2009; 6.8 fb^{-1} = ~8.5 fb^{-1} delivered in 2010 (*if funded...*)

Two benchmark scenarios:

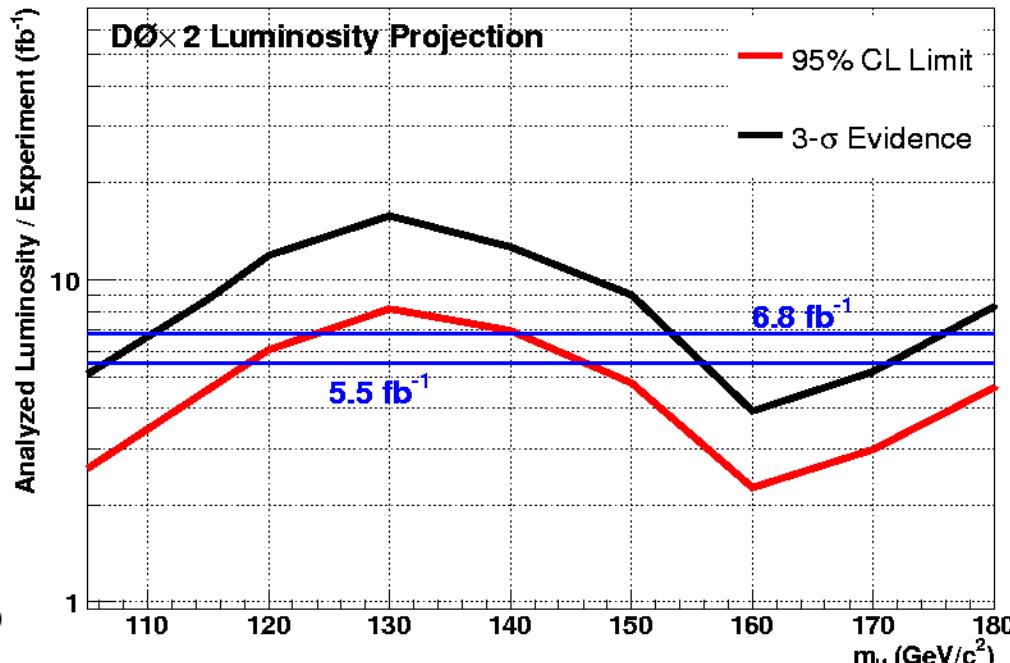
Coasting: no improvements beyond what we have TODAY

Known improvements: achieving potential for known sources of improvement

Coasting



Known Improvements



Sensitivity Projections



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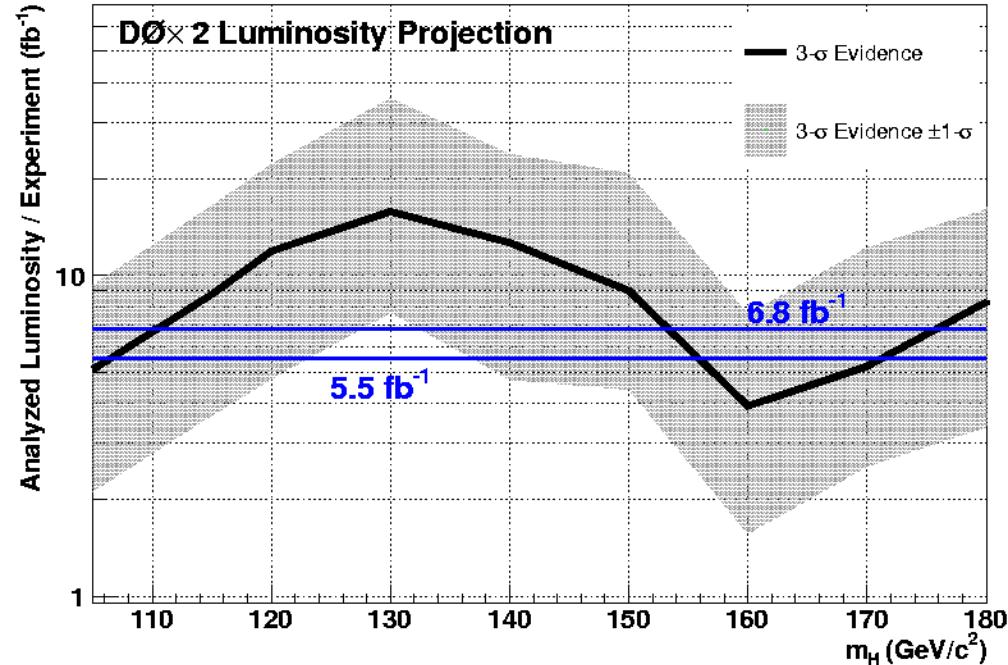
What if we're lucky (or *unlucky*)?

Shaded bands indicate regions enclosing 68% of simulated data outcomes

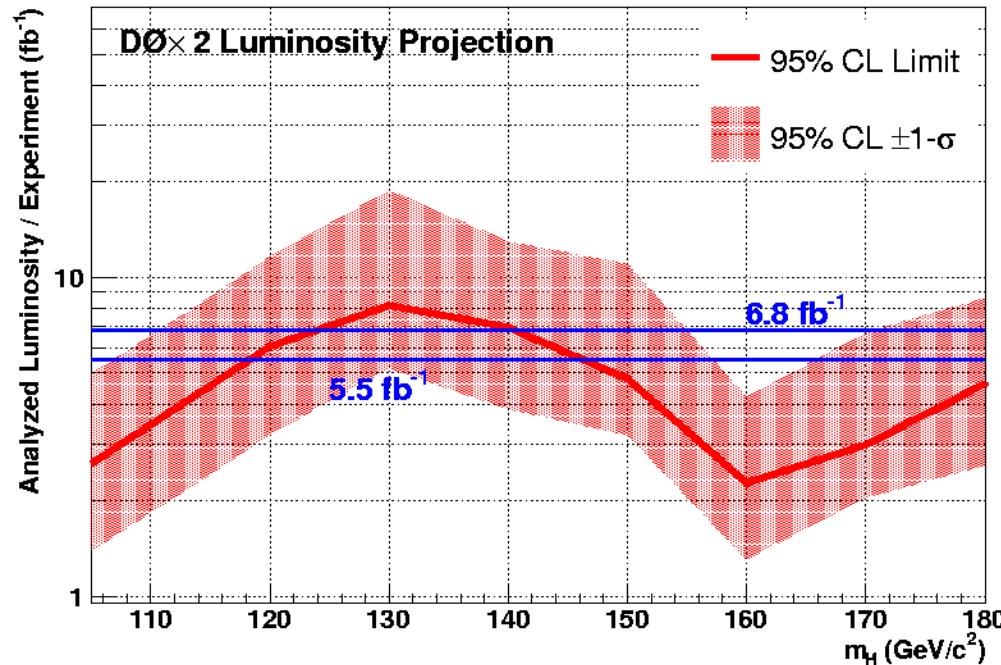
As expected, depends strongly on what we observe!

Region near $m_H=130$ GeV still difficult...

3-Sigma Evidence



95% CL Exclusion



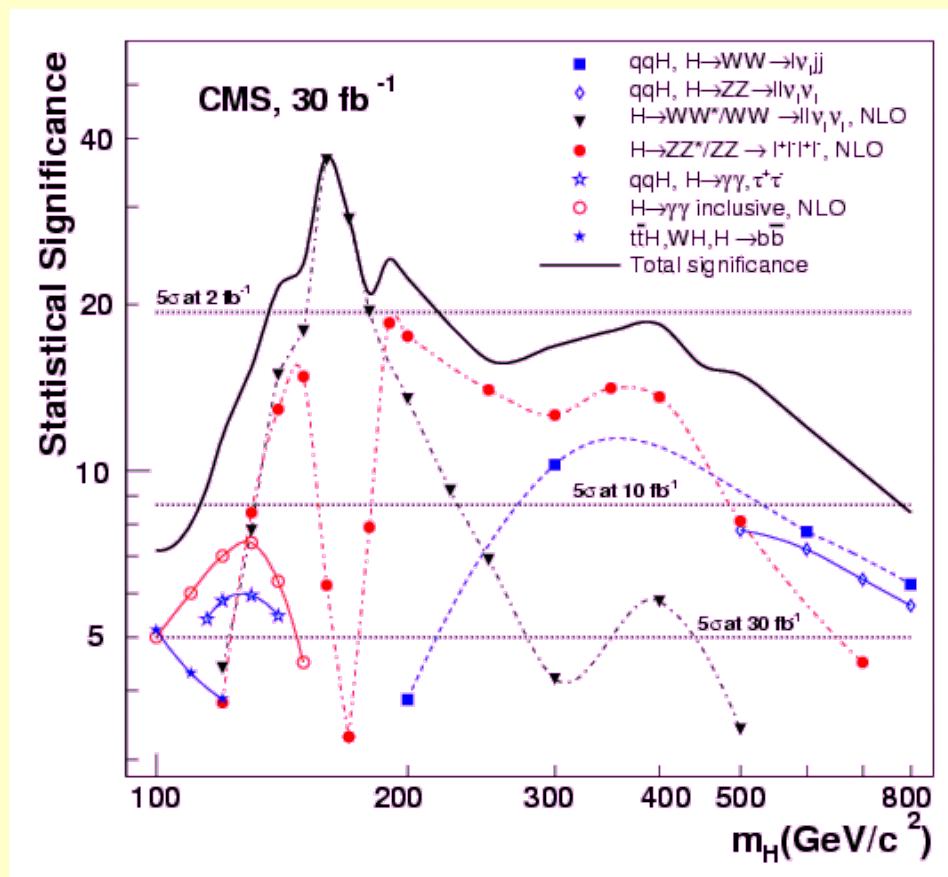
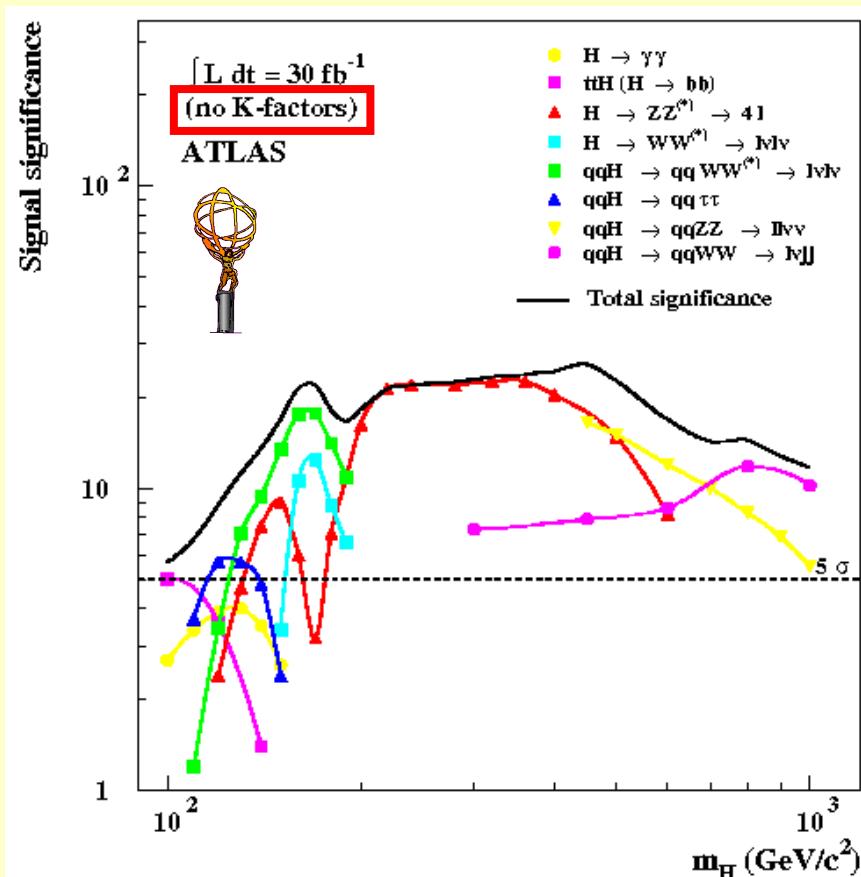
LHC Timescale

LHC predictions show broad discovery potential in 10 fb^{-1} (exclusion requires less)

Low mass Higgs is particularly challenging at LHC

First LHC Higgs search results should be available by summer to fall 2009

Tevatron search will be very mature at this point



Conclusions

- ✗ The Tevatron experiments have demonstrated a robust Higgs program in their first 2 fb^{-1} of recorded data

DZero Only: Observed limits at a ratio of **8.3 (2.5)** to the SM at
 $mH = \mathbf{115 \ (160) \ GeV}$

Tevatron combined: Observed limits at a ratio of **7.8 (1.4)** to the SM at
 $mH = \mathbf{115 \ (160) \ GeV}$

Expect 7-8 fb^{-1} by August 2009, more if 2010 running approved

Working to demonstrate “faster-than-lumi” increases in sensitivity

- ✗ New analyses and improvements frequently available

CDF & DØ will soon update low mass analyses, DØ will update at high mass, CDF will deliver $WH \rightarrow WWW$
- ✗ Higgs physics in Run II of the Tevatron is very exciting

Timing of first data from the LHC will make this an interesting race